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Homogenization of a Dirichlet semilinear elliptic problem with a strong singularity at u = 0in a domain with many small holes



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ABSTRACT

In the present paper we perform the homogenization of the semilinear elliptic problem

$$\begin{cases} u^{\varepsilon} \geq 0 & \text{in } \Omega^{\varepsilon}, \\ -div A(x)Du^{\varepsilon} = F(x, u^{\varepsilon}) & \text{in } \Omega^{\varepsilon}, \\ u^{\varepsilon} = 0 & \text{on } \partial \Omega^{\varepsilon}. \end{cases}$$

In this problem F(x, s) is a Carathéodory function such that $0 \leq F(x, s) \leq h(x)/\Gamma(s)$ a.e. $x \in \Omega$ for every s > 0, with h in some $L^{r}(\Omega)$ and Γ a $C^{1}([0, +\infty[)$ function such that $\Gamma(0) = 0$ and $\Gamma'(s) > 0$ for every s > 0. On the other hand the open sets Ω^{ε} are obtained by removing many small holes from a fixed open set Ω in such a way that a "strange term" μu^{0} appears in the limit equation in the case where the function F(x, s) depends only on x.

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We already treated this problem in the case of a "mild singularity", namely in the case where the function F(x,s) satisfies $0 \leq F(x,s) \leq h(x)(\frac{1}{s}+1)$. In this case the solution u^{ε} to the problem belongs to $H_0^1(\Omega^{\varepsilon})$ and its definition is a "natural" and rather usual one.

In the general case where F(x, s) exhibits a "strong singularity" at u = 0, which is the purpose of the present paper, the solution u^{ε} to the problem only belongs to $H^1_{loc}(\Omega^{\varepsilon})$ but in general does not belong to $H^0_0(\Omega^{\varepsilon})$ anymore, even if u^{ε} vanishes on $\partial \Omega^{\varepsilon}$ in some sense. Therefore we introduced a new notion of solution (in the spirit of the solutions defined by transposition) for problems with a strong singularity. This definition allowed us to obtain existence, stability and uniqueness results.

In the present paper, using this definition, we perform the homogenization of the above semilinear problem and we prove that in the homogenized problem, the "strange term" μu^0 still appears in the left-hand side while the source term $F(x, u^0)$ is not modified in the right-hand side.

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1. Introduction

The present paper deals with the homogenization of the following strongly singular semilinear problem posed in perforated domains Ω^{ε} :

$$\begin{cases} u^{\varepsilon} \ge 0 & \text{in } \Omega^{\varepsilon}, \\ -div A(x)Du^{\varepsilon} = F(x, u^{\varepsilon}) & \text{in } \Omega^{\varepsilon}, \\ u^{\varepsilon} = 0 & \text{on } \partial \Omega^{\varepsilon}. \end{cases}$$
(1.0^{\varepsilon})

Here A(x) is a $N \times N$ bounded coercive matrix, F(x, s) is a Carathéodory function $F(x, s) : \Omega \times [0, +\infty[\rightarrow [0, +\infty]]$ which possibly has a very strong singularity at s = 0; an example of such a function F(x, s) is

$$F(x,s) = f(x)\frac{\left(a+\sin\left(\frac{1}{s}\right)\right)}{\exp\left(-\frac{1}{s}\right)} + g(x)\frac{\left(b+\sin\left(\frac{1}{s}\right)\right)}{s^{\gamma}} + l(x) \text{ a.e. } x \in \Omega, \forall s > 0, \quad (1.1)$$

where $\gamma > 0$, a > 1, b > 1 and where the functions f, g and l are nonnegative; another example is given in (2.5) below. The precise assumptions that we actually make on the function F(x, s) are given in Subsection 2.2 below. On the other hand, the open sets Ω^{ε} are obtained by removing many small closed holes from a fixed open set $\Omega \subset \mathbb{R}^N$, $N \ge 2$. The model example is the case where a bounded open set Ω is perforated by small holes which are closed balls of radius r^{ε} with

$$\begin{cases} r^{\varepsilon} = C_0 \varepsilon^{N/(N-2)} & \text{if } N \ge 3, \\ r^{\varepsilon} = \exp(-C_0/\varepsilon^2) & \text{if } N = 2, \end{cases}$$
(1.2)

which are periodically distributed in \mathbb{R}^N at the vertices of an N-dimensional lattice of cubes of size 2ε . The general framework that we will use for Ω^{ε} in the present paper is (a slight generalization of) the one studied by D. Cioranescu and F. Murat in [3] (see also [15] and [5]); it will be described in details in Subsection 2.3 below.

Note that in (1.0^{ε}) the homogeneous Dirichlet boundary condition is imposed on the whole boundary of Ω^{ε} , which includes the boundary of all the holes. In the "classical" case where the singular semilinear term $F(x, u^{\varepsilon})$ is replaced by a fixed source term $f(x) \in L^2(\Omega)$ which does not depend on u^{ε} , the homogenization in the framework of [3] of problem (1.0^{ε}) leads to a problem where "a strange term" μu^0 appears in the left-hand side, where μ is a bounded nonnegative measure of $H^{-1}(\Omega)$ which depends on the holes and which is actually the asymptotic memory of them.

In [9] we treated the case of problem (1.0^{ε}) where the singularity at s = 0 is mild, namely the case where

$$0 \le F(x,s) \le h(x)(\frac{1}{s^{\gamma}} + 1) \text{ a.e. } x \in \Omega, \forall s > 0, \text{ for some } 0 \le \gamma \le 1.$$
(1.3)

In that paper we proved, for ε fixed, existence, stability and uniqueness results for the solution to (1.0^{ε}) , as well as the homogenization result for perforated domains of the type described above. In this case where (1.3) is satisfied, the solutions to problem (1.0^{ε}) belong to $H_0^1(\Omega^{\varepsilon})$, the equation is intended in the usual weak sense (or more exactly in a slight variant of it), and the test functions that we use belong to $H_0^1(\Omega^{\varepsilon})$.

In contrast, the purpose of the present paper is to treat the case with strong singularities, namely the case where $\gamma > 1$ in (1.3), or more generally where F(x, s) can exhibit any type of singularity at s = 0 (see for example (1.1) above and example (2.5) below). In this case the solutions u^{ε} to (1.0^{ε}) do not in general belong to $H_0^1(\Omega^{\varepsilon})$ (see [14]), but only to $H^1_{\text{loc}}(\Omega^{\varepsilon})$, even if u^{ε} vanishes in some sense on $\partial \Omega^{\varepsilon}$. This induces significant difficulties in order to define a convenient notion of solution, on the first hand in the description of the space which the solution has to belong to, and on the second hand in the definition of the space $\mathcal{V}(\Omega^{\varepsilon})$ of test functions to be used in the equation. In [10] we introduced a notion of solution by defining non standard spaces for the solution and for the test functions, and by writing the equation like in the definition of solutions by transposition introduced by J.-L. Lions and E. Magenes and by G. Stampacchia. This framework, which is recalled in Subsection 3.1 below, allowed us to prove in [10] existence, stability and uniqueness results. The present paper uses this framework and can therefore be considered as a continuation of [10]. It is also a confirmation of the fact that the framework introduced in [10] is robust.

In the present paper we prove that, as in the case studied in [3] where F(x, s) depends only on x, the "strange term" μu^0 appears in the left-hand side of the homogenized problem while the source term $F(x, u^0)$ is not modified in the right-hand side. In other terms, see Theorem 5.1 below, we prove that from any sequence of solutions u^{ε} to (1.0^{ε}) one can extract a subsequence which converges, in a convenient sense, to a solution u^0 to the homogenized problem

$$\begin{cases} u^{0} \geq 0 & \text{in } \Omega, \\ -div A(x)Du^{0} + \mu u^{0} = F(x, u^{0}) & \text{in } \Omega, \\ u^{0} = 0 & \text{on } \partial\Omega. \end{cases}$$
(1.4)

Note that the definition of solution that we use for the solution u^0 to (1.4) is a variant of the definition introduced in [10]. This definition has been introduced in Section 5 of [12] and is recalled in Section 4 below.

This homogenization result was not a priori obvious, since the holes "tend to invade the whole of Ω " (see Remark 6.6 below) and since the source term $F(x, u^{\varepsilon})$ has a singular behaviour at the boundary of the holes.

The method of the proof consists in merging the methods of [3] and of [10]. This however presents some difficulty, since the solution u^{ε} to (1.0^{ε}) in general does not belong to $H_0^1(\Omega^{\varepsilon})$. This leads us (see Section 6 below) to modify the test function w^{ε} used in [3] (which is more or less the difference between 1 and the capacitary potential of the holes in Ω) by introducing a variant z^{ε} of it, which now belongs to the space $\mathcal{V}(\Omega^{\varepsilon})$ of test functions introduced in [10].

In the best of our knowledge, there are only a very few papers concerned with homogenization in the context of this type of singular semilinear problems. In the paper [1] the authors deal with the case where, in a fixed domain Ω , the matrices $A^{\varepsilon}(x)$ wildly vary with ε , remaining uniformly bounded and coercive. For that they use the framework introduced in [2] which is based on the use of strong maximum principle and on the assumption that the function F(x, s) is nonincreasing in s. Note that these properties are never used in the present paper, and neither in [9], [10], [11] and [12]. On the other hand, in the paper [8] the authors consider the homogenization of singular semilinear problems posed in a domain divided in two parts separated by an oscillating interface. Lastly, in the paper [13], the authors study the homogenization in infinite cylinders perforated with small holes with Dirichlet boundary condition. In contrast, there are many papers concerned with existence and uniqueness of solutions to (1.0^{ε}) for ε fixed. Let us just quote, inter alia, [2], [4], [14], [16] and [17].

To conclude this Introduction, let us mention that in the present context of the homogenization of strongly singular semilinear problems posed in perforated domains with Dirichlet boundary condition, we have not been able to prove a corrector result, while we were able to do it in [9] in the context of mild singularities. The corrector result thus remains for us an open problem in the case of strong singularities.

The plan of the present paper is as follows: In Section 2 we give the assumptions that we make on the matrix A(x), on the function F(x, s) and on the sequence of perforated domains Ω^{ε} . In Section 3 we recall the definition introduced in [10] of the solution to the strongly singular semilinear problem posed in Ω^{ε} , and the results of existence, stability and uniqueness obtained in [10]. Note that these solutions satisfy a priori estimates which are recalled in Section 7. In Section 4, we recall the definition given in [12] of the solution to the homogenized problem with a strange term (1.4) (this definition is a variant of the definition given in [10]). In Section 5, we state the main result of the present paper, namely the homogenization result for problem (1.0^{ε}) . This result is proved in Section 8. An important tool for this proof, namely the function z^{ε} which replaces here the function w^{ε} used in [3], is defined in Section 6.

2. Notation and assumptions

As said in the Introduction, in this paper we deal with the asymptotic behaviour, as ε tends to zero, of solutions to the singular semilinear elliptic problem

$$\begin{cases} u^{\varepsilon} \ge 0 & \text{in } \Omega^{\varepsilon}, \\ -div A(x)Du^{\varepsilon} = F(x, u^{\varepsilon}) & \text{in } \Omega^{\varepsilon}, \\ u^{\varepsilon} = 0 & \text{on } \partial\Omega^{\varepsilon}, \end{cases}$$
(2.0^{\varepsilon})

where F(x, s) is possibly singular at s = 0, where u^{ε} satisfies the homogeneous Dirichlet boundary condition on the whole of the boundary of Ω^{ε} , and where Ω^{ε} is a perforated domain obtained by removing many small holes from a given open bounded set Ω in \mathbb{R}^N , with a repartition of those many small holes producing a "strange term" when ε tends to 0.

After the brief Subsection 2.1 dealing with some notation, we give in Subsection 2.2 the assumptions on the matrix A(x) and on the function F(x, s); then in Subsection 2.3 we describe the geometry of the perforated domains and (a slight generalization of) the framework introduced in [3] for treating this problem when the right-hand side is F(x, u) = f(x) in $L^2(\Omega)$.

2.1. Notation

In this paper Ω denotes a bounded open subset of \mathbb{R}^N .

We denote by $\mathcal{D}(\Omega)$ the space of the $C^{\infty}(\Omega)$ functions whose support is compact and included in Ω , and by $\mathcal{D}'(\Omega)$ the space of distributions on Ω .

We denote by $\mathcal{M}_{h}^{+}(\Omega)$ the space of nonnegative bounded Radon measures on Ω .

Since Ω is bounded, $||Dw||_{L^2(\Omega)^N}$ is a norm which is equivalent to $||w||_{H^1(\Omega)}$ on $H^1_0(\Omega)$. We set

$$\|w\|_{H^1_0(\Omega)} = \|Dw\|_{(L^2(\Omega))^N} \quad \forall w \in H^1_0(\Omega).$$

For every $s \in \mathbb{R}$ and every k > 0 we define as usual

$$s^+ = \max\{s, 0\}, \ s^- = \max\{0, -s\},$$

 $T_k(s) = \max\{-k, \min\{s, k\}\}, \ G_k(s) = s - T_k(s).$

For any measurable function $l: x \in \Omega \to l(x) \in [0, +\infty]$ we denote

$$\{l=0\}=\{x\in\Omega: l(x)=0\}, \ \ \{l>0\}=\{x\in\Omega: l(x)>0\}.$$

Finally, in the present paper, we denote by φ functions which belong to $H_0^1(\Omega) \cap L^{\infty}(\Omega)$, while we denote by ϕ functions which belong to $\mathcal{D}(\Omega)$.

2.2. The matrix A(x) and the function F(x, s)

In this Subsection, we give the precise assumptions that we make on the data of problem (2.0^{ϵ}) .

We assume that

$$\Omega$$
 is an open bounded set of \mathbb{R}^N , $N \ge 2$, (2.1)

(no regularity is assumed on the boundary $\partial \Omega$ of Ω), that the matrix A is bounded and coercive, i.e. satisfies

$$A(x) \in (L^{\infty}(\Omega))^{N \times N}, \ \exists \alpha > 0, \ A(x) \ge \alpha I \quad \text{a.e. } x \in \Omega,$$
(2.2)

and that the function F satisfies

$$\begin{aligned} F:(x,s) &\in \Omega \times [0,+\infty[\to F(x,s) \in [0,+\infty] \text{ is a Carathéodory function,} \\ \text{i.e. } F \text{ satisfies} \\ i) \,\forall s \in [0,+\infty[, \, x \in \Omega \to F(x,s) \in [0,+\infty] \text{ is measurable,} \\ ii) \text{ for a.e. } x \in \Omega, \, s \in [0,+\infty[\to F(x,s) \in [0,+\infty] \text{ is continuous,} \end{aligned}$$

$$(2.3)$$

$$\begin{aligned} &(i) \exists h, h(x) \geq 0 \text{ a.e. } x \in \Omega, h \in L^{r}(\Omega), \\ &\text{with } r = \frac{2N}{N+2} \text{ if } N \geq 3, r > 1 \text{ if } N = 2, \\ ⅈ) \exists \Gamma : s \in [0, +\infty[\to \Gamma(s) \in [0, +\infty[, \Gamma \in C^{1}([0, +\infty[), (2.4)])]) \\ &\text{ such that } \Gamma(0) = 0 \text{ and } \Gamma'(s) > 0 \forall s > 0, \\ &iii) 0 \leq F(x, s) \leq \frac{h(x)}{\Gamma(s)} \text{ a.e. } x \in \Omega, \forall s > 0. \end{aligned}$$

Remark 2.1.

• i) Note that in the whole of the present paper we assume that

$$N \ge 2,$$

(see Remark 2.2 below).

- *ii*) Note that the matrix A(x) and the function F(x, s) are defined for $x \in \Omega$ and not only for $x \in \Omega^{\varepsilon}$.
- *iii*) The function F(x, s) can have a very wild behaviour in s when s tends to zero. A possible example is given by (1.1) above, or more generally by

$$F(x,s) = f(x)\frac{(a+\sin(S(s)))}{\exp(-S(s))} + g(x)\frac{(b+\sin(\frac{1}{s}))}{s^{\gamma}} + l(x) \text{ a.e. } x \in \Omega, \, \forall s > 0, \quad (2.5)$$

where $\gamma > 0$, a > 1, b > 1, where the function S satisfies

$$S \in C^{1}(]0, +\infty[), \quad S'(s) < 0 \ \forall s > 0, \quad S(s) \to +\infty \ \text{as} \ s \to 0,$$
 (2.6)

and where the functions f, g and l are nonnegative and belong to $L^{r}(\Omega)$ with r defined by (2.4) above. About example (2.5), please see Remark 2.1 *viii*) of [10].

- *iv*) The function F = F(x, s) is a nonnegative Carathéodory function with values in $[0, +\infty]$ and not only in $[0, +\infty[$. But, in view of conditions (2.4 *ii*) and (2.4 *iii*), for almost every $x \in \Omega$, the function F(x, s) can take the value $+\infty$ only when s = 0 (or, in other terms, F(x, s) is finite for almost every $x \in \Omega$ when s > 0).
- v) Note that the growth condition (2.4 *iii*) is stated for every s > 0, while in (2.3) F is supposed to be a Carathéodory function defined for s in [0, +∞[and not only in]0, +∞[. Indeed an indeterminacy ⁰/₀ appears in ^{h(x)}/_{Γ(s)} when h(x) = 0 and s = 0, while the growth and Carathéodory assumptions (2.4) and (2.3) imply that

$$F(x,s) = 0 \quad \forall s \ge 0 \quad \text{a.e. on } \{x \in \Omega : h(x) = 0\}$$

On the other hand, when h is assumed to satisfy h(x) > 0 for almost every $x \in \Omega$, one can write (2.4 *iii*) for every $s \ge 0$.

• vi) The function h which appears in hypothesis (2.4 i) is an element of $H^{-1}(\Omega)$. Indeed, when $N \ge 3$, the exponent $r = \frac{2N}{N+2}$ is nothing but the Hölder's conjugate $(2^*)'$ of the Sobolev's exponent 2^* , i.e.

when
$$N \ge 3$$
, $\frac{1}{r} = 1 - \frac{1}{2^*}$, where $\frac{1}{2^*} = \frac{1}{2} - \frac{1}{N}$. (2.7)

Making an abuse of notation, we will set

$$2^* = \text{any } p \text{ with } 1 (2.8)$$

With this abuse of notation, h belongs to $L^r(\Omega) = L^{(2^*)'}(\Omega) \subset H^{-1}(\Omega)$ for all $N \ge 2$ since Ω is bounded. This result is indeed a consequence of Sobolev's and Trudinger Moser's inequalities, which (with the above abuse of notation) assert that

$$||v||_{L^{2^*}(\Omega)} \le C_S ||Dv||_{(L^2(\Omega))^N} \quad \forall v \in H^1_0(\Omega) \text{ when } N \ge 2,$$
 (2.9)

where $C_S = C_S(N)$ when $N \ge 3$ and $C_S = C_S(p, \Omega)$ when N = 2. In the latest case, for p given with $1 , the constant <math>C_S = C_S(p, \Omega)$ is bounded independently of Ω when $\Omega \subset Q$, for Q a bounded open set of \mathbb{R}^2 .

• vii) In Section 5 of [9] we performed the homogenization of problem (2.0^{ε}) in the case where F(x, s) has a mild singularity at s = 0, namely in the case where in (2.4 *iii*) the function F(x, s) satisfies

$$0 \le F(x,s) \le h(x)(\frac{1}{s^{\gamma}}+1)$$
 with $0 < \gamma \le 1.$ (2.10)

This is a particular case of the general case treated in the present paper, but that case is easier to treat since the solution u^{ε} to (2.0^{ε}) belongs to $H_0^1(\Omega^{\varepsilon})$ when (2.10)holds true. This property allowed us to prove in that case a corrector result when the matrix A(x) is symmetric and when u^0 further belongs to $L^{\infty}(\Omega)$, see Theorem 5.5 of [9].

Many other remarks can be made about the function F(x, s), and we refer the reader to Section 2 of [10] for them. \Box

2.3. The perforated domains Ω^{ε}

In order to obtain the domain Ω^{ε} , we perforate the fixed domain Ω (see (2.1)) in a way that we describe now. According to (a slight generalization of) the setting presented in [3], we consider here, for every ε which takes its values in a sequence of positive numbers which tends to zero, a finite number $n(\varepsilon)$ of closed sets T_i^{ε} of \mathbb{R}^N , $1 \le i \le n(\varepsilon)$, which are the holes. The domain Ω^{ε} is defined by removing these holes T_i^{ε} from Ω , that is by setting D. Giachetti et al. / Journal of Functional Analysis 274 (2018) 1747-1789

$$\Omega^{\varepsilon} = \Omega - \bigcup_{i=1}^{n(\varepsilon)} T_i^{\varepsilon}.$$
(2.11)

Here, as well as everywhere in the present paper, for every function y^{ε} in $L^{2}(\Omega)$, we define y^{ε} as the extension by zero of y^{ε} to Ω , namely by

$$\widetilde{y^{\varepsilon}}(x) = \begin{cases} y^{\varepsilon}(x) & \text{in } \Omega^{\varepsilon}, \\ & & n^{(\varepsilon)} \\ 0 & \text{in } \bigcup_{i=1}^{n(\varepsilon)} T_i^{\varepsilon}; \end{cases}$$
(2.12)

then $\widetilde{y^{\varepsilon}} \in L^2(\Omega)$ and $\|\widetilde{y^{\varepsilon}}\|_{L^2(\Omega)} = \|y^{\varepsilon}\|_{L^2(\Omega^{\varepsilon})}$; moreover

if
$$y^{\varepsilon} \in H_0^1(\Omega^{\varepsilon})$$
, then $\widetilde{y^{\varepsilon}} \in H_0^1(\Omega)$ with $\widetilde{Dy^{\varepsilon}} = \widetilde{Dy^{\varepsilon}}$ and $\|\widetilde{y^{\varepsilon}}\|_{H_0^1(\Omega)} = \|y^{\varepsilon}\|_{H_0^1(\Omega^{\varepsilon})}$.
(2.13)

We suppose that the sequence of domains Ω^{ε} is such that there exist a sequence of functions w^{ε} , a distribution $\mu \in \mathcal{D}'(\Omega)$ and two sequences of distributions $\mu^{\varepsilon} \in \mathcal{D}'(\Omega)$ and $\lambda^{\varepsilon} \in \mathcal{D}'(\Omega)$ such that¹

$$w^{\varepsilon} \in H^1(\Omega) \cap L^{\infty}(\Omega), \tag{2.14}$$

$$0 \le w^{\varepsilon} \le 1 \text{ a.e. } x \in \Omega, \tag{2.15}$$

$$w^{\varepsilon} = \widetilde{w^{\varepsilon}} \text{ in } \Omega \text{ and } \forall \varphi \in H_0^1(\Omega) \cap L^{\infty}(\Omega), \ w^{\varepsilon} \varphi \in H_0^1(\Omega^{\varepsilon}),$$
 (2.16)

$$w^{\varepsilon} \to 1$$
 in $H^1(\Omega)$ weakly, in $L^{\infty}(\Omega)$ weakly-star and a.e. in Ω as $\varepsilon \to 0$, (2.17)

$$\mu \in H^{-1}(\Omega), \tag{2.18}$$

$$\begin{cases} -div^{t}A(x)Dw^{\varepsilon} = \mu^{\varepsilon} - \lambda^{\varepsilon} \text{ in } \mathcal{D}'(\Omega), \\ \mu^{\varepsilon} \in H^{-1}(\Omega), \lambda^{\varepsilon} \in H^{-1}(\Omega), \\ \mu^{\varepsilon} \ge 0 \text{ in } \mathcal{D}'(\Omega), \\ \mu^{\varepsilon} \to \mu \text{ in } H^{-1}(\Omega) \text{ strongly}, \\ \langle \lambda^{\varepsilon}, \widetilde{y^{\varepsilon}} \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)} = 0 \quad \forall y^{\varepsilon} \in H^{1}_{0}(\Omega^{\varepsilon}). \end{cases}$$

$$(2.19)$$

The model example for Ω^{ε}

The prototype of the examples where assumptions (2.14), (2.15), (2.16), (2.17), (2.18) and (2.19) are satisfied is the case where the matrix A(x) is the identity (and where therefore the operator -div A(x)D is the Laplace's operator $-\Delta$), where $\Omega \subset \mathbb{R}^N$, $N \geq 2$, and where the holes T_i^{ε} are balls of radius r^{ε} with r^{ε} given by

¹ Erratum corrige: please note that at the beginning of (5.4) of [9] we should have added the assertion " $w^{\varepsilon} = \widetilde{w^{\varepsilon}}$ in Ω and" (as it is done in (2.16) in the present paper).

$$\begin{cases} r^{\varepsilon} = C_0 \varepsilon^{N/(N-2)} & \text{if } N \ge 3, \\ \varepsilon^2 \log r^{\varepsilon} \to -C_0 & \text{if } N = 2, \end{cases}$$

for some $C_0 > 0$ (taking $r^{\varepsilon} = \exp(-C_0/\varepsilon^2)$ is the model case for N = 2) which are periodically distributed at the vertices of an N-dimensional lattice of cubes of size 2ε ; in this case the measure μ is given by

$$\begin{cases} \mu = \frac{S_{N-1}(N-2)}{2^N} C_0^{N-2} & \text{if } N \ge 3, \\ \mu = \frac{2\pi}{4} \frac{1}{C_0} & \text{if } N = 2, \end{cases}$$

see e.g. [3] and [15] for more details, and for other examples, in particular for the case where the holes have a different form and/or are distributed on a manifold. \Box

Remark 2.2. In dimension N = 1, there is no sequence w^{ε} which satisfies (2.16) and (2.17) whenever for every ε there exists at least one hole $T_{i_{\varepsilon}}^{\varepsilon}$ with $T_{i_{\varepsilon}}^{\varepsilon} \cap \overline{\Omega} \neq \emptyset$, see Remark 5.1 of [9] for more details. This is the reason why we assume in the present paper that $N \geq 2$. \Box

Some properties of w^{ε} and μ

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One deduces from the first assertion of (2.16) that

$$w^{\varepsilon} = 0$$
 in $\bigcup_{i=1}^{n(\varepsilon)} T_i^{\varepsilon};$ (2.20)

moreover, the second assertion of (2.16) means that for every ε and every $\varphi \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$, there exists a sequence ϕ_n (which depends on ε and on φ) such that

$$\phi_n \in \mathcal{D}(\Omega^{\varepsilon}), \ \widetilde{\phi_n} \to \widetilde{w^{\varepsilon}\varphi} \text{ in } H^1(\Omega).$$
 (2.21)

On the other hand, taking any $\phi^{\varepsilon} \in \mathcal{D}(\Omega^{\varepsilon})$ as test function in the first statement of (2.19) implies that

$$-div^{t}A(x)Dw^{\varepsilon} = \mu^{\varepsilon} \text{ in } \mathcal{D}'(\Omega^{\varepsilon}), \qquad (2.22)$$

which means that the distribution λ^{ε} "only acts on the holes T_i^{ε} ", $i = 1, \dots, n(\varepsilon)$; this fact is also reflected by the last assertion of (2.19).

Taking $v^{\varepsilon} = w^{\varepsilon}\phi$, with $\phi \in \mathcal{D}(\Omega)$, $\phi \ge 0$, as test function in the first assertion of (2.19) we have, thanks to the first assertion of (2.16) and to the last assertion of (2.19),

$$\int_{\Omega} \phi^{t} A(x) Dw^{\varepsilon} Dw^{\varepsilon} + \int_{\Omega} w^{\varepsilon} {}^{t} A(x) Dw^{\varepsilon} D\phi = \langle \mu^{\varepsilon}, w^{\varepsilon} \phi \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)},$$

from which using (2.17) and the fourth statement of (2.19) we deduce that

$$\int_{\Omega} \phi A(x) Dw^{\varepsilon} Dw^{\varepsilon} \to \langle \mu, \phi \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)} \ \forall \phi \in \mathcal{D}(\Omega), \ \phi \ge 0,$$
(2.23)

and therefore using the coercivity (2.2) that

 $\mu \geq 0.$

The distribution $\mu \in H^{-1}(\Omega)$ is therefore also a nonnegative measure. Moreover, using (2.23), (2.2) and (2.17), one deduces that

$$\begin{cases} \forall \phi \in \mathcal{D}(\Omega), \phi \ge 0, \\ \int_{\Omega} \phi \, d\mu = \langle \mu, \phi \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)} = \lim_{\varepsilon} \int_{\Omega} \phi A(x) Dw^{\varepsilon} Dw^{\varepsilon} \le C \|\phi\|_{L^{\infty}(\Omega)}, \end{cases}$$

for a constant C which does not depend on ϕ . Therefore the measure μ is a nonnegative finite Radon measure which satisfies $\int_{\Omega} d\mu \leq C < +\infty$, or in other terms

$$\mu \in \mathcal{M}_b^+(\Omega). \tag{2.24}$$

We will therefore use in the present paper the following (well) known result² (see e.g. [6] Section 1 and [7] Subsection 2.2 for more details):

if $y \in H_0^1(\Omega)$ and if $\nu \in \mathcal{M}_b^+(\Omega) \cap H^{-1}(\Omega)$, then y (or more exactly its quasicontinuous representative for the $H_0^1(\Omega)$ capacity) satisfies

$$\begin{cases} \forall \nu \in \mathcal{M}_b^+(\Omega) \cap H^{-1}(\Omega), \ \forall y \in H_0^1(\Omega), \\ \text{one has } y \in L^1(\Omega; d\nu) \text{ with } \langle \nu, y \rangle_{H^{-1}(\Omega), H_0^1(\Omega)} = \int_{\Omega} y \, d\nu; \end{cases}$$
(2.25)

moreover

$$\begin{cases} \forall \nu \in \mathcal{M}_b^+(\Omega) \cap H^{-1}(\Omega), \ \forall y \in H^1_{loc}(\Omega) \cap L^{\infty}(\Omega), \\ \text{one has } y \in L^{\infty}(\Omega; d\nu) \text{ with } \|y\|_{L^{\infty}(\Omega; d\nu)} = \|y\|_{L^{\infty}(\Omega)}; \end{cases}$$
(2.26)

therefore when $y \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$, then y belongs to $L^1(\Omega; d\nu) \cap L^{\infty}(\Omega; d\nu)$ and thus to $L^p(\Omega; d\nu)$ for every $p, 1 \leq p \leq +\infty$.

² The reader who would not enter in this theory could continue reading the present paper assuming in (2.18) that μ is a function of $L^r(\Omega)$ (with $r = \frac{2N}{N+2}$ if $N \ge 3$ and r > 1 if N = 2) and not only an element of $H^{-1}(\Omega)$.

The limit problem for a source term in $L^2(\Omega)$

When one assumes that the holes T_i^{ε} , $i = 1, \dots, n(\varepsilon)$, are such that the assumptions (2.14), (2.15), (2.16), (2.17), (2.18) and (2.19) hold true, then (see [3], or [15], or [5] for a more general framework), for every $f \in L^2(\Omega)$, the (unique) solution y^{ε} to the linear problem

$$\begin{cases} y^{\varepsilon} \in H_0^1(\Omega^{\varepsilon}), \\ -div A(x)Dy^{\varepsilon} = f \text{ in } \mathcal{D}'(\Omega^{\varepsilon}), \end{cases}$$

$$(2.27)$$

satisfies

$$\widetilde{y^{\varepsilon}} \rightharpoonup y^0$$
 in $H^1_0(\Omega)$,

where y^0 is the (unique) solution to

$$\begin{cases} y^0 \in H^1_0(\Omega) \cap L^2(\Omega; d\mu), \\ -div A(x)Dy^0 + \mu y^0 = f \text{ in } \mathcal{D}'(\Omega), \end{cases}$$

or equivalently to

$$\begin{cases} y^0 \in H_0^1(\Omega) \cap L^2(\Omega; d\mu), \\ \int_{\Omega} A(x) Dy^0 Dz + \int_{\Omega} y^0 z \, d\mu = \int_{\Omega} fz \ \forall z \in H_0^1(\Omega) \cap L^2(\Omega; d\mu). \end{cases}$$
(2.28)

Note that the "strange term" μu_0 which appears in the limit equation (2.28) is the asymptotic memory of the fact that $\widetilde{y^{\varepsilon}}$ was zero on the holes T_i^{ε} .

3. Definition of a solution to the singular semilinear problem in Ω^{ε}

In Subsection 3.1 we first recall the definition of a solution to the singular semilinear problem (2.0^{ε}) which will be used in the present paper; this definition has been introduced in Section 3 of [10]. Then, in Subsection 3.2, we recall the main properties (existence, uniqueness and stability) of such a solution; we will recall in Section 7 below a priori estimates which are satisfied by every such a solution. All these properties have been stated and proved in Sections 4, 5, 6, and 7 of [10].

3.1. The space $\mathcal{V}(\Omega^{\varepsilon})$ of test functions and the definition of a solution to the problem in Ω^{ε}

In order to recall the notion of solution to problem (2.0^{ε}) that we will use in the present paper, we recall the definition of the space $\mathcal{V}(\Omega^{\varepsilon})$ of test functions and a notation (see Section 3 of [10]). **Definition 3.1.** (Definition 3.1 of [10]) The space $\mathcal{V}(\Omega^{\varepsilon})$ is the space of the functions v^{ε} which satisfy

$$v^{\varepsilon} \in H^1_0(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \tag{3.1}$$

$$\begin{cases} \exists I^{\varepsilon} \text{ finite, } \exists \hat{\varphi}_{i}^{\varepsilon}, \exists \hat{g}_{i}^{\varepsilon}, i \in I^{\varepsilon}, \exists \hat{f}^{\varepsilon}, \text{ with} \\ \hat{\varphi}_{i}^{\varepsilon} \in H_{0}^{1}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \hat{g}_{i}^{\varepsilon} \in (L^{2}(\Omega^{\varepsilon}))^{N}, \hat{f}^{\varepsilon} \in L^{1}(\Omega^{\varepsilon}), \\ \text{such that } -div \, {}^{t}\!A(x)Dv^{\varepsilon} = \sum_{i \in I^{\varepsilon}} \hat{\varphi}_{i}^{\varepsilon}(-div \, \hat{g}_{i}^{\varepsilon}) + \hat{f}^{\varepsilon} \text{ in } \mathcal{D}'(\Omega^{\varepsilon}). \end{cases}$$

$$(3.2)$$

In the definition of $\mathcal{V}(\Omega^{\varepsilon})$ we use the notation $\hat{\varphi}_i^{\varepsilon}$, \hat{g}_i^{ε} and \hat{f}^{ε} to help the reader to identify the functions which enter in the definition of the functions of $\mathcal{V}(\Omega^{\varepsilon})$.

Observe that $\mathcal{V}(\Omega^{\varepsilon})$ is a vector space.

Definition 3.2. (Definition 3.2 of [10]) When $v \in \mathcal{V}(\Omega^{\varepsilon})$ with

$$-div \,{}^{t}\!A(x)Dv^{\varepsilon} = \sum_{i\in I^{\varepsilon}} \hat{\varphi}_{i}^{\varepsilon}(-div\,\hat{g}_{i}^{\varepsilon}) + \hat{f}^{\varepsilon} \quad \text{in } \mathcal{D}'(\Omega^{\varepsilon}),$$

where I^{ε} , $\hat{\varphi}_{i}^{\varepsilon}$, $\hat{g}_{i}^{\varepsilon}$ and \hat{f}^{ε} are as in (3.2), and when y^{ε} satisfies

$$y^{\varepsilon} \in H^1_{loc}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}) \text{ with } \varphi^{\varepsilon}y^{\varepsilon} \in H^1_0(\Omega^{\varepsilon}) \, \forall \varphi^{\varepsilon} \in H^1_0(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}),$$

we use the following notation:

$$\langle \langle -div \, {}^{t}\!A(x) Dv^{\varepsilon}, y^{\varepsilon} \rangle \rangle_{\Omega^{\varepsilon}} = \sum_{i \in I^{\varepsilon}} \int_{\Omega^{\varepsilon}} \hat{g}_{i}^{\varepsilon} D(\hat{\varphi}_{i}^{\varepsilon} y^{\varepsilon}) + \int_{\Omega^{\varepsilon}} \hat{f}^{\varepsilon} y^{\varepsilon}. \quad \Box$$
(3.3)

In notation (3.3), the right-hand side is correctly defined since $\hat{\varphi}_i^{\varepsilon} y^{\varepsilon} \in H_0^1(\Omega^{\varepsilon})$ and since $y^{\varepsilon} \in L^{\infty}(\Omega^{\varepsilon})$. In contrast the left-hand side $\langle \langle -div \, {}^t\!A \, Dv^{\varepsilon}, y^{\varepsilon} \rangle \rangle_{\Omega^{\varepsilon}}$ is just a notation.

Remark 3.3. In this Remark we recall some observations which are detailed in Remarks 3.4 and 3.5 of [10].

• i) If $\overline{y^{\varepsilon}} \in H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$, then $\varphi^{\varepsilon} \overline{y^{\varepsilon}} \in H_0^1(\Omega^{\varepsilon})$ for every $\varphi^{\varepsilon} \in H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$, so that for every $v^{\varepsilon} \in \mathcal{V}(\Omega^{\varepsilon})$, $\langle \langle -div \, {}^t\!A \, Dv^{\varepsilon}, y^{\varepsilon} \rangle \rangle_{\Omega^{\varepsilon}}$ is defined. In this case one has

$$\begin{cases} \forall v^{\varepsilon} \in \mathcal{V}(\Omega^{\varepsilon}), \, \forall \overline{y^{\varepsilon}} \in H^{1}_{0}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \\ \langle \langle -div^{t}A(x)Dv^{\varepsilon}, \overline{y^{\varepsilon}} \rangle \rangle_{\Omega^{\varepsilon}} = \langle -div^{t}A(x)Dv^{\varepsilon}, \overline{y^{\varepsilon}} \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})}. \end{cases}$$
(3.4)

• *ii*) If $\varphi^{\varepsilon} \in H^1_0(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$, then $(\varphi^{\varepsilon})^2 \in \mathcal{V}(\Omega^{\varepsilon})$, with

$$-div {}^{t}A(x)D(\varphi^{\varepsilon})^{2} = \hat{\varphi}^{\varepsilon}(-div \,\hat{g}^{\varepsilon}) + \hat{f}^{\varepsilon} \text{ in } \mathcal{D}'(\Omega^{\varepsilon}), \qquad (3.5)$$

with $\hat{\varphi}^{\varepsilon} = 2\varphi^{\varepsilon} \in H^1_0(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \ \hat{g}^{\varepsilon} = {}^tA(x)D\varphi^{\varepsilon} \in (L^2(\Omega^{\varepsilon}))^N$ and $\hat{f}^{\varepsilon} = -2{}^tA(x)D\varphi^{\varepsilon}D\varphi^{\varepsilon} \in L^1(\Omega^{\varepsilon}).$

More in general, if φ_1^{ε} and φ_2^{ε} belong to $H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$, then $\varphi_1^{\varepsilon}\varphi_2^{\varepsilon}$ belongs to $\mathcal{V}(\Omega^{\varepsilon})$.

• *iii*) If $\varphi^{\varepsilon} \in H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$ with supp $\varphi^{\varepsilon} \subset K^{\varepsilon}$, K^{ε} compact, $K^{\varepsilon} \subset \Omega^{\varepsilon}$, then $\varphi^{\varepsilon} \in \mathcal{V}(\Omega^{\varepsilon})$, since

$$-div {}^{t}A(x)D\varphi^{\varepsilon} = \overline{\phi}^{\varepsilon}(-div {}^{t}A(x)D\varphi^{\varepsilon}) \text{ in } \mathcal{D}'(\Omega^{\varepsilon}),$$
(3.6)

for every $\overline{\phi}^{\varepsilon} \in \mathcal{D}(\Omega^{\varepsilon})$, with $\overline{\phi}^{\varepsilon} = 1$ on K^{ε} . In particular every $\phi^{\varepsilon} \in \mathcal{D}(\Omega^{\varepsilon})$ belongs to $\mathcal{V}(\Omega^{\varepsilon})$. \Box

We now recall the definition of a solution to problem (2.0^{ε}) that we will use in the present paper.

Definition 3.4. (Definition 3.6 of [10]) Assume that the matrix A and the function F satisfy (2.2), (2.3) and (2.4). We say that u^{ε} is a solution to problem (2.0^{ε}) if u^{ε} satisfies

$$\begin{cases} i) \ u^{\varepsilon} \in L^{2}(\Omega^{\varepsilon}) \cap H^{1}_{\text{loc}}(\Omega^{\varepsilon}), \\ ii) \ u^{\varepsilon}(x) \geq 0 \quad \text{a.e.} \ x \in \Omega^{\varepsilon}, \\ iii) \ G_{k}(u^{\varepsilon}) \in H^{1}_{0}(\Omega^{\varepsilon}) \quad \forall k > 0, \\ iv) \ \varphi^{\varepsilon}T_{k}(u^{\varepsilon}) \in H^{1}_{0}(\Omega^{\varepsilon}) \quad \forall k > 0, \ \forall \varphi^{\varepsilon} \in H^{1}_{0}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \end{cases}$$
(3.7^{*\varepsilon*})

$$\begin{cases} \forall v^{\varepsilon} \in \mathcal{V}(\Omega^{\varepsilon}), v^{\varepsilon} \geq 0, \\ \text{with } -div^{t}A(x)Dv^{\varepsilon} = \sum_{i \in I^{\varepsilon}} \hat{\varphi}_{i}^{\varepsilon}(-div\,\hat{g}_{i}^{\varepsilon}) + \hat{f}^{\varepsilon} \text{ in } \mathcal{D}'(\Omega^{\varepsilon}), \\ \text{where } \hat{\varphi}_{i}^{\varepsilon} \in H_{0}^{1}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \hat{g}_{i}^{\varepsilon} \in (L^{2}(\Omega^{\varepsilon}))^{N}, \hat{f}^{\varepsilon} \in L^{1}(\Omega^{\varepsilon}), \\ \text{one has} \\ i) \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon})v^{\varepsilon} < +\infty, \\ ii) \int_{\Omega^{\varepsilon}} t^{A}(x)Dv^{\varepsilon}DG_{k}(u^{\varepsilon}) + \sum_{i \in I^{\varepsilon}} \int_{\Omega^{\varepsilon}} \hat{g}_{i}^{\varepsilon}D(\hat{\varphi}_{i}^{\varepsilon}T_{k}(u^{\varepsilon})) + \int_{\Omega^{\varepsilon}} \hat{f}^{\varepsilon}T_{k}(u^{\varepsilon}) = \\ = \langle -div^{t}A(x)Dv^{\varepsilon}, G_{k}(u^{\varepsilon}) \rangle_{H^{-1}(\Omega^{\varepsilon}), H_{0}^{1}(\Omega^{\varepsilon})} + \langle \langle -div^{t}A(x)Dv^{\varepsilon}, T_{k}(u^{\varepsilon}) \rangle \rangle_{\Omega^{\varepsilon}} = \\ = \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon})v^{\varepsilon} \; \forall k > 0. \quad \Box \end{cases}$$

Remark 3.5. When u^{ε} satisfies (3.7^{ε}) , one has

$$\varphi^{\varepsilon} D u^{\varepsilon} \in (L^2(\Omega^{\varepsilon}))^N \quad \forall \varphi^{\varepsilon} \in H^1_0(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon});$$
(3.9)

indeed one writes in $(\mathcal{D}'(\Omega^{\varepsilon}))^N$

$$\varphi^{\varepsilon}Du^{\varepsilon} = \varphi^{\varepsilon}DT_k(u^{\varepsilon}) + \varphi^{\varepsilon}DG_k(u^{\varepsilon}) = D(\varphi^{\varepsilon}T_k(u^{\varepsilon})) - T_k(u^{\varepsilon})D\varphi^{\varepsilon} + \varphi^{\varepsilon}DG_k(u^{\varepsilon}).$$

Remark 3.6. In Definition 3.4, the requirement (3.7^{ε}) is the "space" (which is not a vectorial space) to which the solution should belong, while the requirement $(3.8^{\varepsilon} ii)$ expresses the partial differential equation of (2.0^{ε}) in terms of (non standard) test functions, in the spirit of the solutions defined by transposition introduced by J.-L. Lions and E. Magenes and by G. Stampacchia.

Indeed, very formally, we have

$$\begin{cases} ``\langle -div\,{}^{t}\!A(x)Dv^{\varepsilon},G_{k}(u^{\varepsilon})\rangle_{H^{-1}(\Omega^{\varepsilon}),H^{1}_{0}(\Omega^{\varepsilon})} = \int_{\Omega^{\varepsilon}} (-div\,{}^{t}\!A(x)Dv^{\varepsilon})\,G_{k}(u^{\varepsilon}) = \\ = \int_{\Omega^{\varepsilon}} v^{\varepsilon}\,(-divA(x)DG_{k}(u^{\varepsilon}))", \\ \begin{cases} ``\langle\langle -div\,{}^{t}\!A(x)Dv^{\varepsilon},T_{k}(u^{\varepsilon})\rangle\rangle_{\Omega^{\varepsilon}} = \int_{\Omega^{\varepsilon}} (-div\,{}^{t}\!A(x)Dv^{\varepsilon})\,T_{k}(u^{\varepsilon}) = \\ = \int_{\Omega^{\varepsilon}} v^{\varepsilon}\,(-divA(x)DT_{k}(u^{\varepsilon})", \end{cases}$$

so that $(3.8^{\varepsilon} ii)$ formally means that

$$"\int_{\Omega^{\varepsilon}} v^{\varepsilon} \left(-divA(x)Du^{\varepsilon} \right) = \int_{\Omega^{\varepsilon}} F(x,u^{\varepsilon})v^{\varepsilon} " \ \forall v^{\varepsilon} \in \mathcal{V}(\Omega^{\varepsilon}), \, v^{\varepsilon} \geq 0.$$

Since every v^{ε} can be written as $v^{\varepsilon} = (v^{\varepsilon})^+ - (v^{\varepsilon})^-$ with $(v^{\varepsilon})^+ \ge 0$, $(v^{\varepsilon})^- \ge 0$, one has formally (this is formal since we do not know whether $(v^{\varepsilon})^+$ and $(v^{\varepsilon})^-$ belong to $\mathcal{V}(\Omega^{\varepsilon})$ when v^{ε} belongs to $\mathcal{V}(\Omega^{\varepsilon})$)

$$"-div A(x)Du^{\varepsilon} = F(x, u^{\varepsilon})"$$

which is the second statement of (2.0^{ε}) .

On the other hand, the third assertion of (3.7^{ε}) formally implies (this is formal since in the present paper the boundary $\partial \Omega^{\varepsilon}$ of Ω^{ε} is not assumed to be smooth) that for every k > 0, one has " $G_k(u^{\varepsilon}) = 0$ on $\partial \Omega^{\varepsilon}$ ", i.e. " $0 \le u^{\varepsilon} \le k$ on $\partial \Omega^{\varepsilon}$ " in view of the second assertion of (3.7^{ε}) , which formally implies that " $u^{\varepsilon} = 0$ on $\partial \Omega^{\varepsilon}$ ", i.e. the third statement of (2.0^{ε}) .

For further remarks about Definition 3.4, see Remark 3.7 and Proposition 3.8 of [10]. \Box

3.2. Statements of existence, stability and uniqueness results for the problem in Ω^{ε}

In this Subsection we recall results of existence, stability and uniqueness of the solution to problem (2.0^{ε}) in the sense of Definition 3.4. These results have been stated and proved in [10].

Theorem 3.7. (Existence) (Theorem 4.1 of [10]) Assume that the matrix A and the function F satisfy (2.2), (2.3) and (2.4). Then there exists at least one solution u^{ε} to problem (2.0^{ε}) in the sense of Definition 3.4. \Box

Theorem 3.8. (Stability) (Theorem 4.2 of [10]) Assume that the matrix A satisfies assumption (2.2). Let F_n be a sequence of functions and F_{∞} be a function which all satisfy assumptions (2.3) and (2.4) for the same h and the same Γ . Assume moreover that

a.e.
$$x \in \Omega$$
, $F_n(x, s_n) \to F_\infty(x, s_\infty)$ if $s_n \to s_\infty, s_n \ge 0, s_\infty \ge 0.$ (3.10)

Let u_n^{ε} be any solution to problem $(2.0^{\varepsilon})_n$ in the sense of Definition 3.4, where $(2.0^{\varepsilon})_n$ is the problem (2.0^{ε}) with F(x,s) replaced by $F_n(x,s)$.

Then there exists a subsequence, still labelled by n, and a function u_{∞} , which is a solution to problem $(2.0^{\varepsilon})_{\infty}$ in the sense of Definition 3.4, such that (for ε fixed)

$$\begin{aligned} u_n^{\varepsilon} &\to u_{\infty}^{\varepsilon} \text{ in } L^2(\Omega^{\varepsilon}) \text{ strongly, in } H^1_{loc}(\Omega^{\varepsilon}) \text{ strongly and a.e. in } \Omega^{\varepsilon}, \\ G_k(u_n^{\varepsilon}) &\to G_k(u_{\infty}^{\varepsilon}) \text{ in } H^1_0(\Omega^{\varepsilon}) \text{ strongly } \forall k > 0, \\ \varphi^{\varepsilon} T_k(u_n^{\varepsilon}) &\to \varphi^{\varepsilon} T_k(u_{\infty}^{\varepsilon}) \text{ in } H^1_0(\Omega^{\varepsilon}) \text{ strongly } \forall k > 0, \forall \varphi^{\varepsilon} \in H^1_0(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}). \end{aligned}$$

$$\end{aligned}$$

$$\end{aligned}$$

Finally, the following uniqueness result holds true when, further to (2.3) and (2.4), the function F(x, s) is assumed to be nonincreasing with respect to s, i.e. to satisfy

$$F(x,s) \le F(x,t)$$
 a.e. $x \in \Omega, \ \forall s, \forall t, 0 \le t \le s.$ (3.12)

Theorem 3.9. (Uniqueness) (Theorem 4.3 of [10]) Assume that the matrix A and the function F satisfy (2.2), (2.3) and (2.4). Assume moreover that the function F(x, s) satisfies assumption (3.12). Then the solution to problem (2.0^{ε}) in the sense of Definition 3.4 is unique. \Box

When assumptions (2.2), (2.3), (2.4) as well as (3.12) hold true, Theorems 3.7, 3.8 and 3.9 together assert that problem (2.0^{ε}) is well posed in the sense of Hadamard in the framework of Definition 3.4.

In Section 7 below, we will recall a priori estimates which are satisfied by every solution to (2.0^{ε}) in the sense of Definition 3.4.

4. Definition of a solution to the homogenized singular semilinear problem in Ω

In this Section we recall the definition of the solution to the problem

$$\begin{cases} u \ge 0, & \text{in } \Omega, \\ -div A(x)Du + \mu u = F(x, u) & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$
(4.1)

when μ satisfies

$$\mu \in \mathcal{M}_{h}^{+}(\Omega) \cap H^{-1}(\Omega). \tag{4.2}$$

This Definition, which has been introduced in Section 5 of [12], is an adaptation of Definition 3.4 above.

Definition 4.1. (Definition 5.1 of [12]) Assume that the matrix A, the function F and the Radon measure μ satisfy (2.2), (2.3), (2.4) and (4.2). We say that u is a solution to problem (4.1) if u satisfies

$$\begin{cases} i) \ u \in L^2(\Omega) \cap H^1_{\text{loc}}(\Omega), \\ ii) \ u(x) \ge 0 \text{ a.e. } x \in \Omega, \\ iii) \ G_k(u) \in H^1_0(\Omega) \quad \forall k > 0, \\ iv) \ \varphi T_k(u) \in H^1_0(\Omega) \ \forall k > 0, \ \forall \varphi \in H^1_0(\Omega) \cap L^\infty(\Omega), \end{cases}$$
(4.3)

$$\begin{cases} \forall v \in \mathcal{V}(\Omega), v \geq 0, \\ \text{with } -div^{t}A(x)Dv = \sum_{i \in I} \hat{\varphi}_{i}(-div\,\hat{g}_{i}) + \hat{f} \text{ in } \mathcal{D}'(\Omega), \\ \text{where } \hat{\varphi}_{i} \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega), \hat{g}_{i} \in (L^{2}(\Omega))^{N}, \hat{f}_{i} \in L^{1}(\Omega), \\ \text{one has} \\ i) \int_{\Omega} F(x,u)v < +\infty, \\ ii) \int_{\Omega} F(x,u)v < +\infty, \\ iii) \int_{\Omega} t^{4}A(x)DvDG_{k}(u) + \sum_{i \in I} \int_{\Omega} \hat{g}_{i}D(\hat{\varphi}_{i}T_{k}(u)) + \int_{\Omega} \hat{f}T_{k}(u) + \int_{\Omega} uvd\mu = \\ = \langle -div^{t}A(x)Dv, G_{k}(u) \rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} + \langle \langle -div^{t}A(x)Dv, T_{k}(u) \rangle \rangle_{\Omega} + \int_{\Omega} uv\,d\mu = \\ = \int_{\Omega} F(x,u)\,v \quad \forall k > 0. \quad \Box \end{cases}$$

$$(4.4)$$

Note that the term $\int_{\Omega} uvd\mu$ has a meaning, as shown in the following Remark.

Remark 4.2. In (4.4 *ii*) the term $\int_{\Omega} uvd\mu$ has a meaning since (4.3) and $v \in H_0^1(\Omega) \cap L^\infty(\Omega)$ actually imply that

$$uv \in L^1(\Omega; d\mu). \tag{4.5}$$

Indeed one can write

$$uv = T_k(u)v + G_k(u)v,$$

where $T_k(u)v$ and $G_k(u)$, which belong to $H_0^1(\Omega)$ by (4.3 *iv*) and (4.3 *iii*), belong to $L^1(\Omega; d\mu)$ in view of (2.25), while v, which belongs to $H_0^1(\Omega) \cap L^{\infty}(\Omega)$, belongs to $L^{\infty}(\Omega; d\mu)$ in view of (2.26).

Actually one can prove (see Section 5 of [12]) that every function u which is a solution to problem (4.1) in the sense of Definition 4.1 satisfies the regularity result

$$u \in L^2(\Omega; d\mu). \tag{4.6}$$

Since $v \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$ also belongs to $L^2(\Omega; d\mu)$ in view of (2.25) and (2.26), this again proves (4.5).

Note however that this second proof of (4.5) uses the fact that u satisfies (4.3) and (4.4), while the first proof only uses the fact that u satisfies (4.3). \Box

Remark 4.3. As mentioned in Section 5 of [12], one can prove, for solutions to problem (4.1) in the sense of Definition 4.1, results of existence, stability and uniqueness which are similar to the results recalled in Subsection 3.2 above for the solutions to problem (2.0^{ε}) in the sense of Definition 3.4. Every solution to problem (4.1) in the sense of Definition 4.1 moreover satisfies a priori estimates which are similar to the ones recalled in Section 7 above, see Section 5 of [12] for more details. \Box

5. Statement of the homogenization result for the singular semilinear problem in Ω^{ε}

The existence Theorem 3.7 above asserts that when the matrix A and the function F satisfy assumptions (2.2), (2.3) and (2.4), then for every given $\varepsilon > 0$, the singular semilinear problem (2.0^{ε}) posed in Ω^{ε} has at least a solution u^{ε} in the sense of Definition 3.4; moreover (see Theorem 3.9 above) this solution is unique if the function F(x, s) also satisfies assumption (3.12).

The following result, which is the main result of the present paper, asserts that the homogenization process for the singular semilinear problem (2.0^{ε}) produces a result which is very similar to the homogenization result (2.28) above which holds true for the "classical" problem (2.27) when the source term f belongs to $L^2(\Omega)$.

Theorem 5.1. (Homogenization) Assume that the matrix A and the function F satisfy (2.2), (2.3) and (2.4). Assume also that the sequence of perforated sets Ω^{ε} is such that

(2.14), (2.15), (2.16), (2.17), (2.18) and (2.19) are satisfied. For every $\varepsilon > 0$, let u^{ε} be any solution to problem (2.0^{ε}) in the sense of Definition 3.4, or, in other terms, any function which satisfies (3.7^{ε}) and (3.8^{ε}) , i.e.

$$\begin{cases} i) u^{\varepsilon} \in L^{2}(\Omega^{\varepsilon}) \cap H^{1}_{loc}(\Omega^{\varepsilon}), \\ ii) u^{\varepsilon}(x) \geq 0 \text{ a.e. } x \in \Omega^{\varepsilon}, \\ iii) G_{k}(u^{\varepsilon}) \in H^{1}_{0}(\Omega^{\varepsilon}) \ \forall k > 0, \\ iv) \varphi^{\varepsilon} T_{k}(u^{\varepsilon}) \in H^{1}_{0}(\Omega^{\varepsilon}) \ \forall k > 0, \ \forall \varphi^{\varepsilon} \in H^{1}_{0}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \end{cases}$$
(5.1)

$$\begin{cases} \forall v^{\varepsilon} \in \mathcal{V}(\Omega^{\varepsilon}), v^{\varepsilon} \geq 0, \\ \text{with } - div^{t}A(x)Dv^{\varepsilon} = \sum_{i \in I^{\varepsilon}} \hat{\varphi}_{i}^{\varepsilon}(-div\,\hat{g}_{i}^{\varepsilon}) + \hat{f}^{\varepsilon} \text{ in } \mathcal{D}'(\Omega^{\varepsilon}), \\ \text{where } \hat{\varphi}_{i}^{\varepsilon} \in H_{0}^{1}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \hat{g}_{i}^{\varepsilon} \in L^{2}(\Omega^{\varepsilon})^{N}, \hat{f}^{\varepsilon} \in L^{1}(\Omega^{\varepsilon}), \\ \text{one has} \\ i) \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon})v^{\varepsilon} < +\infty, \\ ii) \int_{\Omega^{\varepsilon}} tA(x)Dv^{\varepsilon}DG_{k}(u^{\varepsilon}) + \sum_{i \in I^{\varepsilon}} \int_{\Omega^{\varepsilon}} \hat{g}_{i}^{\varepsilon}D(\hat{\varphi}_{i}^{\varepsilon}T_{k}(u^{\varepsilon})) + \int_{\Omega} \hat{f}^{\varepsilon}T_{k}(u^{\varepsilon}) = \\ = \langle -div^{t}A(x)Dv^{\varepsilon}, G_{k}(u^{\varepsilon}) \rangle_{H^{-1}(\Omega^{\varepsilon}), H_{0}^{1}(\Omega^{\varepsilon})} + \langle \langle -div^{t}A(x)Dv^{\varepsilon}, T_{k}(u^{\varepsilon}) \rangle \rangle_{\Omega^{\varepsilon}} = \\ = \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon})v^{\varepsilon} \quad \forall k > 0. \end{cases}$$

$$(5.2)$$

Then there exists a subsequence, still denoted by ε , such that for $\widetilde{u^{\varepsilon}}$, the extension by zero of u^{ε} to Ω defined by (2.12), one has

$$\widetilde{u^{\varepsilon}} \rightharpoonup u^0$$
 in $L^2(\Omega)$ weakly and a.e. in Ω , (5.3)

$$G_k(\widetilde{u^{\varepsilon}}) \rightharpoonup G_k(u^0)$$
 in $H^1_0(\Omega)$ weakly and a.e. in $\Omega \ \forall k > 0,$ (5.4)

$$\begin{cases} \varphi w^{\varepsilon} T_k(\widetilde{u^{\varepsilon}}) \rightharpoonup \varphi T_k(u^0) \text{ in } H_0^1(\Omega) \text{ weakly and a.e. in } \Omega\\ \forall k > 0, \forall \varphi \in H_0^1(\Omega) \cap L^{\infty}(\Omega), \end{cases}$$
(5.5)

where u^0 satisfies (4.3) and (4.4), or, in other terms, where the limit u^0 is a solution to problem (4.1) in the sense of Definition 4.1, i.e.

$$\begin{cases} i) u^{0} \in L^{2}(\Omega) \cap H^{1}_{loc}(\Omega), \\ ii) u^{0}(x) \geq 0 \quad \text{a.e. } x \in \Omega, \\ iii) G_{k}(u^{0}) \in H^{1}_{0}(\Omega) \quad \forall k > 0, \\ iv) \varphi T_{k}(u^{0}) \in H^{1}_{0}(\Omega) \quad \forall k > 0, \quad \forall \varphi \in H^{1}_{0}(\Omega) \cap L^{\infty}(\Omega), \end{cases}$$
(5.6)

$$\begin{cases} \forall v \in \mathcal{V}(\Omega), v \geq 0, \\ \text{with } -div^{t}A(x)Dv = \sum_{i \in I} \hat{\varphi}_{i}(-div\,\hat{g}_{i}) + \hat{f} \text{ in } \mathcal{D}'(\Omega), \\ \text{where } \hat{\varphi}_{i} \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega), \hat{g}_{i} \in L^{2}(\Omega)^{N}, \hat{f} \in L^{1}(\Omega), \\ \text{one has} \\ i) \int_{\Omega} F(x, u^{0})v < +\infty, \\ (5.7) \\ ii) \int_{\Omega} t^{A}(x)DvDG_{k}(u^{0}) + \sum_{i \in I} \int_{\Omega} \hat{g}_{i}D(\hat{\varphi}_{i}T_{k}(u^{0})) + \int_{\Omega} \hat{f}T_{k}(u^{0}) + \int_{\Omega} uv \, d\mu = \\ = \langle -div^{t}A(x)Dv, G_{k}(u^{0}) \rangle_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} + \langle \langle -div^{t}A(x)Dv, T_{k}(u^{0}) \rangle \rangle_{\Omega} + \int_{\Omega} uv \, d\mu = \\ = \int_{\Omega} F(x, u^{0}) \, \forall k > 0. \quad \Box \end{cases}$$

Remark 5.2. Observe that in the case where assumption (3.12) is made on the function F(x, s), i.e. when F(x, s) is assumed to be nonincreasing with respect to s, the solutions u^{ε} to problem (2.0^{ε}) in the sense of Definition 3.4 and u^0 to problem (4.1) in the sense of Definition 4.1 are unique. In this case there is no need to extract a subsequence in Theorem 5.1 and the convergences (5.3), (5.4) and (5.5) hold true for the whole sequence ε . \Box

6. Definition of the function z^{ε} , and the strong convergence of $\chi_{\Omega^{\varepsilon}}$

6.1. Definition of the function z^{ε} , a variant of the test function w^{ε}

The idea of the proof of the homogenization Theorem 5.1 of the present paper is to combine the ideas of the proof of the existence Theorem 4.1 of [10] with the ideas of the proof of the homogenization Theorem 1.2 of [3]. In the latest paper a key tool is the use of the test function $w^{\varepsilon}\phi$, where $\phi \in \mathcal{D}(\Omega)$ and where w^{ε} is defined is in (2.14), (2.15), (2.16), (2.17), (2.18) and (2.19). Unfortunately, this function does not (seem to) belong to $\mathcal{V}(\Omega^{\varepsilon})$: indeed, the function $w^{\varepsilon}\phi$ belongs to $H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$, but the computation in $\mathcal{D}'(\Omega^{\varepsilon})$ of $-div {}^tA(x)D(w^{\varepsilon}\phi)$ produces four terms, where three of them are in the form required for $w^{\varepsilon}\phi$ to belong to $\mathcal{V}(\Omega^{\varepsilon})$, but where the fourth term

$$\phi(-div^{t}ADw^{\varepsilon}) = \phi(\mu^{\varepsilon} - \lambda^{\varepsilon}) = \phi\mu^{\varepsilon} \text{ in } H^{-1}(\Omega^{\varepsilon}),$$

is in the form $\phi(-\operatorname{div} G^{\varepsilon})$ for some $G^{\varepsilon} \in (L^2(\Omega^{\varepsilon}))^N$ (see (6.10) below), but not in the required form $\hat{\varphi}_i^{\varepsilon}(-\operatorname{div} \hat{G}^{\varepsilon})$ with $\hat{\varphi}_i^{\varepsilon} \in H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$, since ϕ belongs to $H_0^1(\Omega) \cap L^{\infty}(\Omega)$ but not to $H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$. For this reason we introduce in this Section the function z^{ε} , which is a variant of w^{ε} but which is such that $z^{\varepsilon}v$ belongs to $\mathcal{V}(\Omega^{\varepsilon})$ for every $v \in \mathcal{V}(\Omega)$. Note that the function z^{ε} does not (seem to) belong to the smaller (and

easier to understand) space of test functions $\mathcal{W}(\Omega^{\varepsilon})$ introduced in Subsection 3.3 of [12], which is made of finite sums of products $\varphi^{\varepsilon}\psi^{\varepsilon}$ with φ^{ε} and ψ^{ε} in $H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$. This is actually the reason for which we decided to choose in [10] the framework of the space $\mathcal{V}(\Omega)$ instead of the framework of the space $\mathcal{W}(\Omega)$.

Proposition 6.1. Assume that (2.14), (2.15), (2.16), (2.17), (2.18) and (2.19) hold true. Then (for a subsequence, as far as the almost everywhere convergence in (6.5) is concerned), there exists a function z^{ε} such that

$$z^{\varepsilon} \in H^1(\Omega) \cap L^{\infty}(\Omega), \tag{6.1}$$

$$z^{\varepsilon} - w^{\varepsilon} \in H^1_0(\Omega^{\varepsilon}), \tag{6.2}$$

$$0 \le z^{\varepsilon}(x) \le w^{\varepsilon}(x) \le 1 \quad \text{a.e. } x \in \Omega, \tag{6.3}$$

$$z^{\varepsilon}v \in \mathcal{V}(\Omega^{\varepsilon}) \quad \forall v \in \mathcal{V}(\Omega),$$
 (6.4)

$$z^{\varepsilon} \to 1$$
 in $H^1(\Omega)$ weakly, in $L^{\infty}(\Omega)$ weakly-star and a.e. in Ω as $\varepsilon \to 0$, (6.5)

$$-div {}^{t}A(x)Dz^{\varepsilon} = w^{\varepsilon}\mu^{\varepsilon} \quad \text{in } \mathcal{D}'(\Omega^{\varepsilon}). \quad \Box$$
(6.6)

Remark 6.2. Note that in view of (2.20), the two first assertions of (6.3) imply that in particular

$$z^{\varepsilon} = 0 \text{ in } \bigcup_{i=1}^{n(\varepsilon)} T_i^{\varepsilon}. \quad \Box$$
(6.7)

Remark 6.3. As far as (6.4) is concerned, we will actually prove that

$$z^{\varepsilon}\varphi \in H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}) \quad \forall \varphi \in H_0^1(\Omega) \cap L^{\infty}(\Omega),$$
(6.8)

and that if v is such that

$$\begin{cases} v \in \mathcal{V}(\Omega) \\ \text{with } -div \,{}^{t}\!A(x)Dv = \sum_{i \in I} \hat{\varphi}_{i}(-div\,\hat{g}_{i}) + \hat{f} \text{ in } \mathcal{D}'(\Omega), \\ \text{where } \hat{\varphi}_{i} \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega), \hat{g}_{i} \in L^{2}(\Omega)^{N}, \hat{f} \in L^{1}(\Omega), \end{cases}$$
(6.9)

and if G^{ε} is a sequence such that

$$\begin{cases} \mu^{\varepsilon} = -div \, G^{\varepsilon} \text{ in } \mathcal{D}'(\Omega), \ \mu = -div \, G \text{ in } \mathcal{D}'(\Omega), \\ \text{with } G^{\varepsilon} \to G \text{ in } (L^2(\Omega))^N \text{ strongly,} \end{cases}$$
(6.10)

(note that such a sequence G^{ε} exists since μ^{ε} converges strongly in $H^{-1}(\Omega)$ to μ in view of the fourth assertion of (2.19)), one has

$$\begin{cases} z^{\varepsilon}v \in \mathcal{V}(\Omega^{\varepsilon}) \\ \text{with } -div \,{}^{t}\!A(x)D(z^{\varepsilon}v) = \\ = \sum_{i \in I} z^{\varepsilon}\hat{\varphi}_{i} \left(-div \,\hat{g}_{i}\right) + w^{\varepsilon}v \left(-div \,G^{\varepsilon}\right) + z^{\varepsilon}\hat{f} - {}^{t}\!A(x)DvDz^{\varepsilon} - {}^{t}\!A(x)Dz^{\varepsilon}Dv \right. \tag{6.11}$$

in $\mathcal{D}'(\Omega^{\varepsilon})$. \Box

Proof of Proposition 6.1.

First step. Since $w^{\varepsilon} \in H^1(\Omega)$ and $\mu^{\varepsilon} \in H^{-1}(\Omega)$, the product $w^{\varepsilon}\mu^{\varepsilon}$ is, as usual, the distribution on Ω defined, for every $\phi \in \mathcal{D}(\Omega)$, as

$$\langle w^{\varepsilon} \mu^{\varepsilon}, \phi \rangle_{\mathcal{D}'(\Omega), \mathcal{D}(\Omega)} = \langle \mu^{\varepsilon}, w^{\varepsilon} \phi \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)}.$$
(6.12)

We claim that actually

$$w^{\varepsilon}\mu^{\varepsilon} \in H^{-1}(\Omega); \tag{6.13}$$

indeed, since $\mu^{\varepsilon} \geq 0$ in $\mathcal{D}'(\Omega)$ (see the third assertion of (2.19)), μ^{ε} is a nonnegative Radon measure on Ω , and therefore μ^{ε} belongs to $\mathcal{M}_{b}^{+}(\omega)$ for every open set ω with $\overline{\omega} \subset \Omega$. Taking, for any given $\phi \in \mathcal{D}(\Omega)$, an open set ω with $\operatorname{supp} \phi \subset \omega \subset \overline{\omega} \subset \Omega$, we have, using (6.12) in Ω and then (2.25) in ω ,

$$\langle w^{\varepsilon}\mu^{\varepsilon},\phi\rangle_{\mathcal{D}'(\Omega),\mathcal{D}(\Omega)} = \langle \mu^{\varepsilon},w^{\varepsilon}\phi\rangle_{H^{-1}(\Omega),H^{1}_{0}(\Omega)} = \langle \mu^{\varepsilon},w^{\varepsilon}\phi\rangle_{H^{-1}(\omega),H^{1}_{0}(\omega)} = \int\limits_{\omega} w^{\varepsilon}\phi\,d\mu^{\varepsilon},$$

and then, using (2.15) and then (2.25) in Ω ,

$$\begin{cases} |\langle w^{\varepsilon}\mu^{\varepsilon},\phi\rangle_{\mathcal{D}'(\Omega),\mathcal{D}(\Omega)}| = \left|\int_{\omega} w^{\varepsilon}\phi\,d\mu^{\varepsilon}\right| \leq \int_{\omega} w^{\varepsilon}|\phi|d\mu^{\varepsilon} \leq \int_{\omega} |\phi|\,d\mu^{\varepsilon} = \\ = \int_{\Omega} |\phi|\,d\mu^{\varepsilon} = \langle \mu^{\varepsilon},|\phi|\rangle_{H^{-1}(\Omega),H^{1}_{0}(\Omega)} \leq \|\mu^{\varepsilon}\|_{H^{-1}(\Omega)}\|\phi\|_{H^{1}_{0}(\Omega)} \,\,\forall\phi\in\mathcal{D}(\Omega). \end{cases}$$

This implies that (6.13) holds true with

$$\|w^{\varepsilon}\mu^{\varepsilon}\|_{H^{-1}(\Omega)} \le \|\mu^{\varepsilon}\|_{H^{-1}(\Omega)}.$$
(6.14)

Second step. Since $w^{\varepsilon} \mu^{\varepsilon}$ belongs to $H^{-1}(\Omega)$ by (6.13), one has

$$w^{\varepsilon}\mu^{\varepsilon} \in H^{-1}(\Omega) \subset H^{-1}(\Omega^{\varepsilon}).$$

Applying Lax–Milgram's lemma then implies the existence (and the uniqueness) of the solution y^{ε} to

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$$\begin{cases} y^{\varepsilon} \in H^{1}(\Omega^{\varepsilon}), \\ y^{\varepsilon} - w^{\varepsilon} \in H^{1}_{0}(\Omega^{\varepsilon}), \\ -div^{t}A(x)Dy^{\varepsilon} = w^{\varepsilon}\mu^{\varepsilon} \quad \text{in } \mathcal{D}'(\Omega^{\varepsilon}). \end{cases}$$
(6.15)

We now define z^{ε} by

$$z^{\varepsilon} = \widetilde{y^{\varepsilon}}, \tag{6.16}$$

where $\widetilde{y^{\varepsilon}}$ is the extension by zero of $\underline{y^{\varepsilon}}$ to Ω defined by (2.12). Therefore z^{ε} satisfies (6.2) and (6.6), and one deduces from $\widetilde{y^{\varepsilon} - w^{\varepsilon}} = z^{\varepsilon} - \widetilde{w^{\varepsilon}}$, from $\widetilde{y^{\varepsilon} - w^{\varepsilon}} \in H_0^1(\Omega)$, from the first assertion of (2.16) and from (2.14) that $z^{\varepsilon} \in H^1(\Omega)$.

Third step. We now prove that

$$0 \le z^{\varepsilon}(x) \le w^{\varepsilon}(x) \text{ a.e. } x \in \Omega, \tag{6.17}$$

a fact which in particular implies (6.3) in view of (2.15), and which completes the proof of (6.1).

Since one has

$$0 = z^{\varepsilon} = w^{\varepsilon} \text{ in } \bigcup_{i=1}^{n(\varepsilon)} T_i^{\varepsilon},$$

a fact which follows from (6.16) and from the first assertion of (2.16), we only have to prove that

$$0 \le y^{\varepsilon}(x) \le w^{\varepsilon}(x)$$
 a.e. $x \in \Omega^{\varepsilon}$. (6.18)

In order to prove (6.18), we first observe that $-(y^{\varepsilon})^{-} \in H_0^1(\Omega^{\varepsilon})$: indeed one has $-(y^{\varepsilon})^{-} \in H^1(\Omega^{\varepsilon})$ in view of (6.15) and also

$$-(y^{\varepsilon}-w^{\varepsilon})^{-} \leq -(y^{\varepsilon})^{-} \leq 0$$
 a.e. in Ω^{ε} ,

where the first inequality results from the facts that the function $-s^-$ is nondecreasing and that $w^{\varepsilon} \ge 0$ (see (2.15)); therefore Lemma A.1 of [10] and $y^{\varepsilon} - w^{\varepsilon} \in H_0^1(\Omega^{\varepsilon})$ imply that $-(y^{\varepsilon})^- \in H_0^1(\Omega^{\varepsilon})$. Using $-(y^{\varepsilon})^-$ as test function in (6.15) we get, in view of (2.15) and of the third assertion of (2.19),

$$\int_{\Omega^{\varepsilon}} {}^{t} A(x) Dy^{\varepsilon} D(-(y^{\varepsilon})^{-}) = \langle w^{\varepsilon} \mu^{\varepsilon}, -(y^{\varepsilon})^{-} \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})} \le 0,$$

which implies that

$$0 \leq y^{\varepsilon}(x)$$
 a.e. $x \in \Omega^{\varepsilon}$.

On the other hand, since $y^{\varepsilon} - w^{\varepsilon} \in H_0^1(\Omega^{\varepsilon})$ by (6.15), using $(y^{\varepsilon} - w^{\varepsilon})^+ \in H_0^1(\Omega^{\varepsilon})$ as test function in (6.15) and (2.22) we get

$$\int_{\Omega^{\varepsilon}} {}^{t} A(x) D(y^{\varepsilon} - w^{\varepsilon}) D(y^{\varepsilon} - w^{\varepsilon})^{+} = \langle w^{\varepsilon} \mu^{\varepsilon} - \mu^{\varepsilon}, (y^{\varepsilon} - w^{\varepsilon})^{+} \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})}$$

Since in view of (2.15) and of the third assertion of (2.19) one has

$$\langle (w^{\varepsilon}-1)\mu^{\varepsilon}, (y^{\varepsilon}-w^{\varepsilon})^{+} \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})} \leq 0,$$

this implies that

$$y^{\varepsilon}(x) - w^{\varepsilon}(x) \le 0$$
 a.e. $x \in \Omega^{\varepsilon}$.

We have proved that (6.18) (and therefore (6.17)) holds true.

Fourth step. Let us now prove that

$$z^{\varepsilon} - w^{\varepsilon} \to 0$$
 in $H_0^1(\Omega)$ strongly. (6.19)

Combined with (2.17) this will imply (6.5) (for a subsequence, as far as the almost everywhere convergence is concerned).

Using $y^{\varepsilon} - w^{\varepsilon} \in H_0^1(\Omega^{\varepsilon})$ as test function in (6.15) and (2.22) we get

$$\int_{\Omega^{\varepsilon}} {}^{t} A(x) D(y^{\varepsilon} - w^{\varepsilon}) D(y^{\varepsilon} - w^{\varepsilon}) = \langle w^{\varepsilon} \mu^{\varepsilon} - \mu^{\varepsilon}, y^{\varepsilon} - w^{\varepsilon} \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})}.$$

By the coercivity of the matrix A, this implies that

$$\begin{cases} \alpha \| y^{\varepsilon} - w^{\varepsilon} \|_{H_0^1(\Omega^{\varepsilon})}^2 \leq \langle w^{\varepsilon} \mu^{\varepsilon} - \mu^{\varepsilon}, y^{\varepsilon} - w^{\varepsilon} \rangle_{H^{-1}(\Omega^{\varepsilon}), H_0^1(\Omega^{\varepsilon})} \leq \\ \leq \left(\| w^{\varepsilon} \mu^{\varepsilon} \|_{H^{-1}(\Omega^{\varepsilon})} + \| \mu^{\varepsilon} \|_{H^{-1}(\Omega^{\varepsilon})} \right) \| y^{\varepsilon} - w^{\varepsilon} \|_{H_0^1(\Omega^{\varepsilon})}, \end{cases}$$

$$(6.20)$$

which in view of (6.14), of the fourth assertion of (2.19) and of (2.13) implies that $\|y^{\varepsilon} - w^{\varepsilon}\|_{H_0^1(\Omega^{\varepsilon})} = \|z^{\varepsilon} - w^{\varepsilon}\|_{H_0^1(\Omega)}$ is bounded. But $z^{\varepsilon} - w^{\varepsilon}$ is also bounded in $L^{\infty}(\Omega)$ in view of (6.3) and (2.15). Therefore $(w^{\varepsilon} - 1)(z^{\varepsilon} - w^{\varepsilon})$ is bounded in $H_0^1(\Omega) \cap L^{\infty}(\Omega)$, so that in view of (2.17)

$$(w^{\varepsilon} - 1)(z^{\varepsilon} - w^{\varepsilon}) \rightarrow 0$$
 in $H_0^1(\Omega)$ weakly. (6.21)

Writing

$$\begin{cases} \langle w^{\varepsilon}\mu^{\varepsilon} - \mu^{\varepsilon}, y^{\varepsilon} - w^{\varepsilon} \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})} = \langle w^{\varepsilon}\mu^{\varepsilon} - \mu^{\varepsilon}, z^{\varepsilon} - w^{\varepsilon} \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)} = \\ = \langle \mu^{\varepsilon}, (w^{\varepsilon} - 1)(z^{\varepsilon} - w^{\varepsilon}) \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)}, \end{cases}$$

and using (6.21) and the fact that μ^{ε} tends to μ in $H^{-1}(\Omega)$ strongly by the fourth assertion of (2.19), we deduce (6.19) from the first line of (6.20).

Fifth step. At this point, we have proved the existence of a sequence z^{ε} which satisfies (6.1), (6.2), (6.3), (6.5) and (6.6). Let us now prove that z^{ε} satisfies (6.4), or more precisely (6.8) and (6.11) when $v \in \mathcal{V}(\Omega)$ satisfies (6.9).

Assertion (6.8) follows from the equality

$$z^{\varepsilon}\varphi = (z^{\varepsilon} - w^{\varepsilon})\varphi + w^{\varepsilon}\varphi,$$

and from the facts that when $\varphi \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$, then both $(z^{\varepsilon} - w^{\varepsilon})\varphi$ and $w^{\varepsilon}\varphi$ belong to $H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$ (see (6.2), (6.3), (2.14), (2.15) and (2.16)).

On the other hand, using (6.9), (6.15) and (6.10), we have in $\mathcal{D}'(\Omega^{\varepsilon})$

$$\begin{cases} -div^{t}A(x)D(z^{\varepsilon}v) = -div(z^{\varepsilon}^{t}A(x)Dv) - div(v^{t}A(x)Dz^{\varepsilon}) = \\ = z^{\varepsilon}(-div^{t}A(x)Dv) - {}^{t}A(x)DvDz^{\varepsilon} + \\ + v(-div^{t}A(x)Dz^{\varepsilon}) - {}^{t}A(x)Dz^{\varepsilon}Dv = \\ = \sum_{i \in I} z^{\varepsilon}\hat{\varphi}_{i}(-div\,\hat{g}_{i}) + z^{\varepsilon}\hat{f} - {}^{t}A(x)DvDz^{\varepsilon} + \\ + vw^{\varepsilon}(-div\,G^{\varepsilon}) - {}^{t}A(x)Dz^{\varepsilon}Dv \text{ in } \mathcal{D}'(\Omega^{\varepsilon}), \end{cases}$$

$$(6.22)$$

which proves that the equation stated in (6.11) holds true. This proves (6.11), and therefore (6.4), in particular since $z^{\varepsilon}\hat{\varphi}_i$ and $w^{\varepsilon}v$ belong to $H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$, by (6.8), (2.16) and (2.14). \Box

6.2. Strong convergence of the sequence $\chi_{\Omega^{\varepsilon}}$ in $L^{1}(\Omega)$

In this Subsection we prove the following Proposition:

Proposition 6.4. Assume that the sequence of perforated set Ω^{ε} is such that (2.14), (2.15), (2.16), (2.17), (2.18) and (2.19) are satisfied. Then

$$\chi_{\Omega^{\varepsilon}} \to 1 \text{ in } L^1(\Omega) \text{ strongly as } \varepsilon \to 0. \quad \Box$$
 (6.23)

From (6.23) one immediately deduces that for a subsequence, still denoted by ε , one has

$$\chi_{\Omega^{\varepsilon}} \to 1 \text{ a.e. in } \Omega \text{ as } \varepsilon \to 0.$$
 (6.24)

Proof. In view of (2.20) one has

$$w^{\varepsilon}\chi_{\Omega^{\varepsilon}} = w^{\varepsilon}$$
 a.e. in Ω . (6.25)

Since $0 \leq \chi_{\Omega^{\varepsilon}} \leq 1$, one can extract a subsequence such that

$$\chi_{\Omega^{\varepsilon}} \rightharpoonup \theta \text{ in } L^{\infty}(\Omega) \text{ weakly-star as } \varepsilon \to 0,$$
(6.26)

so that using (2.17) and passing to the limit in (6.25), one has

$$\theta = 1$$
 a.e. in Ω ,

which implies that (6.26) holds true with $\theta = 1$ for the whole sequence ε .

It is then sufficient to write that

$$\|\chi_{\Omega^{\varepsilon}} - 1\|_{L^{1}(\Omega)} = \int_{\Omega} |\chi_{\Omega^{\varepsilon}} - 1| = \int_{\Omega} (1 - \chi_{\Omega^{\varepsilon}}) \to \int_{\Omega} (1 - \theta) = 0 \text{ as } \varepsilon \to 0$$

to deduce (6.23) for the whole sequence ε . \Box

Remark 6.5. Note that (6.24) implies that, for every subsequence ε' of ε and for almost every $x_0 \in \Omega$, there exists $\varepsilon_0(x_0) > 0$ such that

$$\chi_{\Omega^{\varepsilon'}}(x_0) = 1 \ \forall \varepsilon', \ \varepsilon' < \varepsilon_0(x_0),$$

or in other terms that

$$x_0 \in \Omega^{\varepsilon'} \ \forall \varepsilon', \ \varepsilon' < \varepsilon_0(x_0). \quad \Box \tag{6.27}$$

Remark 6.6. Assertion (6.27) implies that almost every $x_0 \in \Omega$ belongs to $\Omega^{\varepsilon'}$ for ε' sufficiently small ($\varepsilon' < \varepsilon_0(x_0)$), which formally means that " Ω^{ε} is very close to Ω ".

In contrast, note that if we consider the case of holes periodically distributed at the vertices of a cubic lattice of size $\varepsilon_j = 1/2^j$, with $j \in \mathbb{N}$, namely the case considered in the model example described in the Section 2 above, every point of the form

$$c_k^{\varepsilon_{j_0}} = k \,\varepsilon_{j_0} = \left(\frac{k_1}{2^{j_0}}, \frac{k_2}{2^{j_0}}, \cdots, \frac{k_N}{2^{j_0}}\right) \text{ with } k = (k_1, k_2, \cdots, k_N) \in \mathbb{Z}^N \text{ and } j_0 \in \mathbb{N}$$

is, for every $j \geq j_0$, the center of some hole $T_i^{\varepsilon_j}$ which is extremely small, since its size is $r^{\varepsilon_j} = C_0 \varepsilon_j^{N/(N-2)} = C_0 2^{-jN/(N-2)}$ (for $N \geq 3$); therefore such a point $c_k^{\varepsilon_{j_0}}$ does not belong to Ω^{ε} for $\varepsilon_j = 1/2^j \leq 1/2^{j_0}$; note that these points are dense in Ω , and "tend to invade the whole of Ω " as ε_j tends to zero. \Box

7. A priori estimates for the solutions to the singular semilinear problem in Ω^{ε}

In this Section we recall a priori estimates which are satisfied by every solution to (2.0^{ε}) in the sense of Definition 3.4; these estimates have been stated and proved in [10].

Proposition 7.1. (A priori estimate of $G_k(u^{\varepsilon})$ in $H_0^1(\Omega^{\varepsilon})$) (Proposition 5.1 of [10]) Assume that the matrix A and the function F satisfy (2.2), (2.3) and (2.4). Then for every u^{ε} solution to problem (2.0^{ε}) in the sense of Definition 3.4 one has

$$\|G_k(u^{\varepsilon})\|_{H^1_0(\Omega^{\varepsilon})} = \|DG_k(u^{\varepsilon})\|_{(L^2(\Omega^{\varepsilon}))^N} \le \frac{C_S}{\alpha} \frac{\|h\|_{L^r(\Omega^{\varepsilon})}}{\Gamma(k)} \ \forall k > 0, \tag{7.1}$$

where C_S is the (generalized) Sobolev's constant defined by (2.9) above. \Box

Remark 7.2. (Remark 5.2 of [10]) From Poincaré's inequality

$$\|y^{\varepsilon}\|_{L^{2}(\Omega^{\varepsilon})} \leq C_{P}(\Omega^{\varepsilon}) \|Dy^{\varepsilon}\|_{(L^{2}(\Omega^{\varepsilon}))^{N}} \quad \forall y^{\varepsilon} \in H^{1}_{0}(\Omega^{\varepsilon}),$$
(7.2)

where the constant $C_P(\Omega^{\varepsilon})$ is bounded independently of Ω^{ε} when $\Omega^{\varepsilon} \subset Q$, for Q a bounded open set of \mathbb{R}^N , one deduces from (7.1), writing $u^{\varepsilon} = T_k(u^{\varepsilon}) + G_k(u^{\varepsilon})$, that every solution u^{ε} to problem (2.0^{ε}) in the sense of Definition 3.4 satisfies the following a priori estimate in $L^2(\Omega^{\varepsilon})$

$$\|u^{\varepsilon}\|_{L^{2}(\Omega^{\varepsilon})} \leq k|\Omega^{\varepsilon}|^{\frac{1}{2}} + C_{P}(\Omega^{\varepsilon})\frac{C_{S}}{\alpha}\frac{\|h\|_{L^{r}(\Omega^{\varepsilon})}}{\Gamma(k)} \ \forall k > 0,$$

$$(7.3)$$

which, taking $k = k_0$ for some k_0 fixed or minimizing in k provides an a priori estimate of $\|u^{\varepsilon}\|_{L^2(\Omega^{\varepsilon})}$ which does not depend on k. \Box

Proposition 7.3. (A priori estimate of $\varphi^{\varepsilon}DT_k(u^{\varepsilon})$ in $(L^2(\Omega^{\varepsilon}))^N$ for $\varphi^{\varepsilon} \in H_0^1(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon})$) (Proposition 5.4 of [10]) Assume that the matrix A and the function F satisfy (2.2), (2.3) and (2.4). Then for every u^{ε} solution to problem (2.0^{ε}) in the sense of Definition 3.4 one has

$$\begin{cases} \|\varphi^{\varepsilon} DT_{k}(u^{\varepsilon})\|_{(L^{2}(\Omega^{\varepsilon}))^{N}}^{2} \leq \\ \leq \frac{32k^{2}}{\alpha^{2}} \|A\|_{(L^{\infty}(\Omega^{\varepsilon}))^{N\times N}}^{2} \|D\varphi^{\varepsilon}\|_{(L^{2}(\Omega^{\varepsilon}))^{N}}^{2} + \frac{C_{S}^{2}}{\alpha^{2}} \frac{\|h\|_{L^{r}(\Omega^{\varepsilon})}^{2}}{\Gamma(k)^{2}} \|\varphi^{\varepsilon}\|_{L^{\infty}(\Omega^{\varepsilon})}^{2} \\ \forall k > 0, \ \forall \varphi^{\varepsilon} \in H_{0}^{1}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \end{cases}$$
(7.4)

where C_S is the (generalized) Sobolev's constant defined by (2.9) above. \Box

Remark 7.4. (Remark 5.5 of [10]) From the a priori estimate (7.4) one deduces that every solution u^{ε} to problem (2.0^{ε}) in the sense of Definition 3.4 satisfies the following a priori estimate of $\varphi^{\varepsilon}T_k(u^{\varepsilon})$ in $H_0^1(\Omega^{\varepsilon})$

$$\begin{cases} \|\varphi^{\varepsilon}T_{k}(u^{\varepsilon})\|_{H_{0}^{1}(\Omega^{\varepsilon})}^{2} = \|D(\varphi^{\varepsilon}T_{k}(u^{\varepsilon}))\|_{(L^{2}(\Omega^{\varepsilon}))^{N}}^{2} \leq \\ \leq \left(\frac{64k^{2}}{\alpha^{2}}\|A\|_{(L^{\infty}(\Omega^{\varepsilon}))^{N\times N}}^{2} + 2k^{2}\right)\|D\varphi^{\varepsilon}\|_{(L^{2}(\Omega^{\varepsilon}))^{N}}^{2} + 2\frac{C_{S}^{2}}{\alpha^{2}}\frac{\|h\|_{L^{r}(\Omega^{\varepsilon})}^{2}}{\Gamma(k)^{2}}\|\varphi^{\varepsilon}\|_{L^{\infty}(\Omega^{\varepsilon})}^{2} (7.5) \\ \forall k > 0, \ \forall \varphi^{\varepsilon} \in H_{0}^{1}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}). \quad \Box \end{cases}$$

For $\delta > 0$, define the function $Z_{\delta} : s \in [0, +\infty[\to Z_{\delta}(s) \in [0, +\infty[$ by

$$Z_{\delta}(s) = \begin{cases} 1 & \text{if } 0 \le s \le \delta, \\ -\frac{s}{\delta} + 2 & \text{if } \delta \le s \le 2\delta, \\ 0 & \text{if } 2\delta \le s. \end{cases}$$
(7.6)

Proposition 7.5. (Control of the quantity $\int_{\Omega} F(x, u^{\varepsilon}) Z_{\delta}(u^{\varepsilon}) v$ when δ is small) (Proposi-

tion 5.9 of [10]) Assume that the matrix A and the function F satisfy (2.2), (2.3) and (2.4). Then for every u^{ε} solution to problem (2.0^{ε}) in the sense of Definition 3.4 and for every v^{ε} such that

$$\begin{cases} v^{\varepsilon} \in \mathcal{V}(\Omega^{\varepsilon}), v^{\varepsilon} \ge 0, \\ \text{with } -div^{t}A(x)Dv^{\varepsilon} = \sum_{i \in I^{\varepsilon}} \hat{\varphi}_{i}^{\varepsilon}(-div\,\hat{g}_{i}^{\varepsilon}) + \hat{f}^{\varepsilon} \text{ in } \mathcal{D}'(\Omega^{\varepsilon}) \\ \text{where } \hat{\varphi}_{i}^{\varepsilon} \in H_{0}^{1}(\Omega^{\varepsilon}) \cap L^{\infty}(\Omega^{\varepsilon}), \hat{g}_{i}^{\varepsilon} \in L^{2}(\Omega^{\varepsilon})^{N}, \hat{f}^{\varepsilon} \in L^{1}(\Omega^{\varepsilon}), \end{cases}$$
(7.7)

one has

$$\begin{cases} \forall \delta > 0, \ \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) Z_{\delta}(u^{\varepsilon}) v^{\varepsilon} \leq \\ \leq \frac{3}{2} \left(\int_{\Omega^{\varepsilon}} \left| \sum_{i \in I^{\varepsilon}} \hat{g}_{i}^{\varepsilon} D \hat{\varphi}_{i}^{\varepsilon} + \hat{f}^{\varepsilon} \right| \right) \delta + \sum_{i \in I^{\varepsilon}} \int_{\Omega} Z_{\delta}(u^{\varepsilon}) \hat{g}_{i}^{\varepsilon} D u^{\varepsilon} \hat{\varphi}_{i}^{\varepsilon}. \quad \Box \end{cases}$$

$$(7.8)$$

Note that the second term of the right-hand side of (7.8) has a meaning since in view of (3.9) above one has $Du^{\varepsilon} \hat{\varphi}_i^{\varepsilon} \in (L^2(\Omega^{\varepsilon}))^N$.

A consequence of Proposition 7.5 is:

Proposition 7.6. (Proposition 5.12 of [10]) Assume that the matrix A and the function F satisfy (2.2), (2.3) and (2.4). Then for every u^{ε} solution to problem (2.0^{ε}) in the sense of Definition 3.4 one has

$$\int_{\{u^{\varepsilon}=0\}} F(x, u^{\varepsilon})v^{\varepsilon} = 0 \quad \forall v^{\varepsilon} \in \mathcal{V}(\Omega^{\varepsilon}), v^{\varepsilon} \ge 0. \quad \Box$$

$$(7.9)$$

8. Proof of the homogenization Theorem 5.1

First step. In this first step we state a priori estimates and we extract a subsequence, still denoted by ε , such that convergences (5.3), (5.4) and (5.5) of Theorem 5.1 hold true for some u^0 which satisfies (5.6).

As already said in the existence Theorem 3.7 of Subsection 3.2 above, for every ε there exists at least one solution to problem (2.0^{ε}) in the sense of Definition 3.4. This solution in particular satisfies the a priori estimates (7.1), (7.3) and (7.5) stated in Proposition 7.1 and in Remarks 7.2 and 7.4 above.

Since $\Omega^{\varepsilon} \subset \Omega$, since the generalized Sobolev's constant C_S which appears in (2.9) does not depend on Ω^{ε} when $N \geq 3$, and is bounded independently of Ω^{ε} when N = 2since $\Omega^{\varepsilon} \subset \Omega$ (see the comment after (2.9)), and since the Poincaré's constant $C_P(\Omega^{\varepsilon})$ which appears in (7.2) is bounded independently of Ω^{ε} since $\Omega^{\varepsilon} \subset \Omega$ (see the comment after (7.2)), the a priori estimates (7.1) and (7.3) imply that

$$\|G_k(\widetilde{u^{\varepsilon}})\|_{H^1_0(\Omega)} = \|G_k(u^{\varepsilon})\|_{H^1_0(\Omega^{\varepsilon})} \le C(k),$$

$$(8.1)$$

$$\|u^{\varepsilon}\|_{L^{2}(\Omega)} = \|u^{\varepsilon}\|_{L^{2}(\Omega^{\varepsilon})} \le C,$$
(8.2)

where the constants C(k) and C do not depend on ε for k > 0 fixed.

Similarly, taking in (7.5) $\varphi^{\varepsilon} = z^{\varepsilon}\varphi$, with $\varphi \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$ and z^{ε} defined by Proposition 6.1, and observing that $\|z^{\varepsilon}\varphi\|_{L^{\infty}(\Omega^{\varepsilon})}$ and $\|D(z^{\varepsilon}\varphi)\|_{(L^2(\Omega))^N}$ are bounded independently of ε , one obtains that

$$\|z^{\varepsilon}\varphi T_{k}(\widetilde{u^{\varepsilon}})\|_{H^{1}_{0}(\Omega)} = \|z^{\varepsilon}\varphi T_{k}(u^{\varepsilon})\|_{H^{1}_{0}(\Omega^{\varepsilon})} \le C(k)(\|\varphi\|_{L^{\infty}(\Omega)} + \|D\varphi\|_{(L^{2}(\Omega))^{N}}), \quad (8.3)$$

where the constant C(k) does not depend on ε for k > 0 fixed.

Using the (generalized) Sobolev's inequality (2.9) for $G_k(\widetilde{u^{\varepsilon}})$, the fact that z^{ε} is bounded in $H^1(\Omega) \cap L^{\infty}(\Omega)$, $\varphi \in H^1_0(\Omega) \cap L^{\infty}(\Omega)$ and (8.1), one obtains that

$$\begin{cases} \|z^{\varepsilon}\varphi G_{k}(\widetilde{u^{\varepsilon}})\|_{W_{0}^{1,q}(\Omega)} \leq C(k)(\|\varphi\|_{L^{\infty}(\Omega)} + \|D\varphi\|_{(L^{2}(\Omega))^{N}}) \\ \text{where } q \text{ is defined by } \frac{1}{q} = \frac{1}{2^{*}} + \frac{1}{2}, \end{cases}$$

$$(8.4)$$

where 2^* is defined by (2.7) and (2.8), and where the constant C(k) does not depend on ε for k > 0 fixed.

Collecting together (8.3) and (8.4) implies that

$$z^{\varepsilon}\varphi \widetilde{u^{\varepsilon}} = z^{\varepsilon}\varphi T_k(\widetilde{u^{\varepsilon}}) + z^{\varepsilon}\varphi G_k(\widetilde{u^{\varepsilon}})$$

is bounded in $W_0^{1,q}(\Omega)$, and therefore that

$$z^{\varepsilon} \varphi \widetilde{u^{\varepsilon}}$$
 is compact in $L^{q}(\Omega)$ for every $\varphi \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega)$. (8.5)

On the other hand, let us write, for every $\varphi \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$

$$\varphi \widetilde{u^{\varepsilon}} = z^{\varepsilon} \varphi \widetilde{u^{\varepsilon}} + (1 - z^{\varepsilon}) \varphi \widetilde{u^{\varepsilon}}.$$
(8.6)

Since $(1 - z^{\varepsilon})$ tends to zero in $L^{p}(\Omega)$ strongly for every $p < +\infty$ (see (6.5)), since $\varphi \in L^{\infty}(\Omega)$ and since u^{ε} is bounded in $L^{2}(\Omega)$ (see (8.2)), one has

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$$(1 - z^{\varepsilon})\varphi \widetilde{u^{\varepsilon}} \to 0 \text{ in } L^{q}(\Omega) \text{ strongly as } \varepsilon \to 0.$$
 (8.7)

From (8.5), (8.6) and (8.7) one concludes that

$$\varphi \widetilde{u^{\varepsilon}}$$
 is compact in $L^{q}(\Omega) \quad \forall \varphi \in H^{1}_{0}(\Omega) \cap L^{\infty}(\Omega).$ (8.8)

In view of (8.2) and (8.8) one can extract a subsequence, still denoted by ε , such that there exists some $u^0 \in L^2(\Omega)$ such that

$$\widetilde{u^{\varepsilon}} \rightharpoonup u^0$$
 in $L^2(\Omega)$ weakly and a.e. in Ω as $\varepsilon \to 0$. (8.9)

This proves (5.3).

For the same subsequence, one has, in view of (8.9), (8.1), (6.5) and (8.3)

$$G_{k}(\widetilde{u^{\varepsilon}}) \rightharpoonup G_{k}(u^{0}) \text{ in } H^{1}_{0}(\Omega) \text{ weakly and a.e. in } \Omega \quad \forall k > 0 \text{ as } \varepsilon \to 0,$$

$$\begin{cases} z^{\varepsilon} \varphi T_{k}(\widetilde{u^{\varepsilon}}) \rightharpoonup \varphi T_{k}(u^{0}) \text{ in } H^{1}_{0}(\Omega) \text{ weakly and a.e. in } \Omega \\ \forall k > 0, \forall \varphi \in H^{1}_{0}(\Omega) \cap L^{\infty}(\Omega) \text{ as } \varepsilon \to 0. \end{cases}$$

$$(8.11)$$

This proves
$$(5.4)$$
.

Moreover, since similarly to (8.3), one has, taking $\varphi^{\varepsilon} = w^{\varepsilon} \varphi$ in (7.5),

$$\|w^{\varepsilon}\varphi T_{k}(\widetilde{u^{\varepsilon}})\|_{H^{1}_{0}(\Omega)} \leq C(k)(\|\varphi\|_{L^{\infty}(\Omega)} + \|D\varphi\|_{(L^{2}(\Omega))^{N}}),$$
(8.12)

where the constant C(k) does not depend on ε for k > 0 fixed, one also has

$$\begin{cases} w^{\varepsilon}\varphi T_{k}(\widetilde{u^{\varepsilon}}) \rightharpoonup \varphi T_{k}(u^{0}) \text{ in } H_{0}^{1}(\Omega) \text{ weakly and a.e. in } \Omega\\ \forall k > 0, \ \forall \varphi \in H_{0}^{1}(\Omega) \cap L^{\infty}(\Omega) \text{ as } \varepsilon \to 0. \end{cases}$$

$$(8.13)$$

This proves (5.5).

Note that since $\widetilde{u^{\varepsilon}}$ is nonnegative on Ω , one has

$$u^0(x) \ge 0 \text{ a.e. } x \in \Omega. \tag{8.14}$$

On the other hand, since $G_k(u^0) \in H^1_0(\Omega)$ in view of (8.10) and since for every $\phi \in \mathcal{D}(\Omega)$ one has $\phi T_k(u^0) \in H^1_0(\Omega)$ in view of (8.11), the function u^0 satisfies

$$u^0 \in H^1_{\text{loc}}(\Omega). \tag{8.15}$$

As announced at the beginning of this first step, we have extracted a subsequence, still denoted by ε , and defined an u^0 which satisfies (5.6) uch that convergences (5.3), (5.4) and (5.5) hold true.

It remains to prove that u^0 satisfies (5.7). This is done in the second, third, fourth, fifth, sixth, seventh and eighth steps below.

Second step. We now consider any fixed $v \in \mathcal{V}(\Omega)$ with $v \ge 0$. In view of (6.3) and (6.4) and of Remark 6.3, the function $z^{\varepsilon}v$ belongs to $\mathcal{V}(\Omega^{\varepsilon})$ with $z^{\varepsilon}v \ge 0$ and satisfies (6.1) when v satisfies (6.9). The use of $v^{\varepsilon} = z^{\varepsilon}v$ in (5.2) is therefore licit and one has

$$\begin{cases} \int_{\Omega^{\varepsilon}} {}^{t}A(x)D(z^{\varepsilon}v)DG_{k}(u^{\varepsilon}) + \\ + \sum_{i\in I}\int_{\Omega^{\varepsilon}} \hat{g}_{i}D(z^{\varepsilon}\hat{\varphi}_{i}T_{k}(u^{\varepsilon})) + \int_{\Omega^{\varepsilon}} G^{\varepsilon}D(w^{\varepsilon}vT_{k}(u^{\varepsilon})) + \\ + \int_{\Omega^{\varepsilon}} \left(z^{\varepsilon}\hat{f} - {}^{t}A(x)DvDz^{\varepsilon} - {}^{t}A(x)Dz^{\varepsilon}Dv \right) T_{k}(u^{\varepsilon}) = \\ = \langle -div \, {}^{t}A(x)D(z^{\varepsilon}v), G_{k}(u^{\varepsilon}) \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})} + \langle \langle -div \, {}^{t}A(x)D(z^{\varepsilon}v), T_{k}(u^{\varepsilon}) \rangle \rangle_{\Omega^{\varepsilon}} = \\ = \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) z^{\varepsilon}v. \end{cases}$$

$$(8.16)$$

From now on, $v \in \mathcal{V}(\Omega)$, $v \ge 0$, and k > 0 will be fixed.

In the present step and in the next one, we pass to the limit, as ε tends to zero, in the first term of the left-hand side of (8.16) and we prove that

$$\begin{cases} \int_{\Omega^{\varepsilon}} {}^{t} A(x) D(z^{\varepsilon} v) DG_{k}(u^{\varepsilon}) \to \\ \to \int_{\Omega} {}^{t} A(x) Dv DG_{k}(u^{0}) + \int_{\Omega} G_{k}(u^{0}) v \, d\mu \text{ as } \varepsilon \to 0, \end{cases}$$

$$(8.17)$$

in which we observe that the first term of the right-hand side can also be written as

$$\int_{\Omega} {}^{t}A(x)DvDG_{k}(u^{0}) = \langle -div {}^{t}A(x)Dv, G_{k}(u^{0}) \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)}.$$

$$(8.18)$$

For that we introduce, for k > 0 fixed and for every n > k, the function $S_{k,n} : \mathbb{R}^+ \to \mathbb{R}^+$ defined by

$$S_{k,n}(s) = \begin{cases} 0 & \text{if } 0 \le s \le k, \\ s - k & \text{if } k \le s \le n, \\ n - k & \text{if } n \le s. \end{cases}$$
(8.19)

Observe that one has

$$G_k(s) = S_{k,n}(s) + G_n(s) \quad \forall s > 0, \ \forall n, \ n > k,$$
(8.20)

$$S_{k,n}(s) = T_{n-k}(G_k(s)) \quad \forall s > 0, \ \forall n, \ n > k.$$
 (8.21)

Using (8.20) we write the first term of the left-hand side of (8.16) as

$$\begin{cases} \int_{\Omega^{\varepsilon}} {}^{t} A(x) D(z^{\varepsilon} v) DG_{k}(u^{\varepsilon}) = \\ = \int_{\Omega^{\varepsilon}} {}^{t} A(x) D(z^{\varepsilon} v) DS_{k,n}(u^{\varepsilon}) + \int_{\Omega^{\varepsilon}} {}^{t} A(x) D(z^{\varepsilon} v) DG_{n}(u^{\varepsilon}). \end{cases}$$

$$(8.22)$$

We first pass to the limit in the first term of the right-hand side of (8.22) as ε tends to zero for n and k fixed, n > k > 0. For that we write, using (6.6) in the latest equality below,

$$\begin{cases} \int_{\Omega^{\varepsilon}} {}^{t}A(x)D(z^{\varepsilon}v)DS_{k,n}(u^{\varepsilon}) = \\ = \int_{\Omega^{\varepsilon}} {}^{t}A(x)Dz^{\varepsilon}D(S_{k,n}(u^{\varepsilon})v) - \int_{\Omega^{\varepsilon}} {}^{t}A(x)Dz^{\varepsilon}Dv S_{k,n}(u^{\varepsilon}) + \\ + \int_{\Omega^{\varepsilon}} {}^{t}A(x)Dv DS_{k,n}(u^{\varepsilon}) z^{\varepsilon} = \\ = \langle w^{\varepsilon}\mu^{\varepsilon}, S_{k,n}(u^{\varepsilon})v \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})} - \int_{\Omega^{\varepsilon}} {}^{t}A(x)Dz^{\varepsilon}Dv S_{k,n}(u^{\varepsilon}) + \\ + \int_{\Omega^{\varepsilon}} {}^{t}A(x)Dv DS_{k,n}(u^{\varepsilon}) z^{\varepsilon}. \end{cases}$$

$$(8.23)$$

We now observe that in view of formula (8.21) and of the convergence (8.10) of $G_k(\widetilde{u^{\varepsilon}})$ to $G_k(u^0)$ in $H_0^1(\Omega)$ weakly, one has for n > k fixed

$$\begin{cases} S_{k,n}(\widetilde{u^{\varepsilon}}) = T_{n-k}(G_k(\widetilde{u^{\varepsilon}})) \to T_{n-k}(G_k(u^0)) = S_{k,n}(u^0) \\ \text{in } H^1_0(\Omega) \text{ weakly, in } L^{\infty}(\Omega) \text{ weakly-star and a.e. in } \Omega \text{ as } \varepsilon \to 0. \end{cases}$$
(8.24)

Therefore, using in the first term of the right-hand side of (8.23) the strong convergence of μ^{ε} to μ in $H^{-1}(\Omega)$ (see the fourth assertion of (2.19)) and the convergence (2.17), and then the equality (2.25), we have

$$\begin{cases} \langle w^{\varepsilon} \mu^{\varepsilon}, S_{k,n}(u^{\varepsilon}) v \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})} = \langle \mu^{\varepsilon}, S_{k,n}(u^{\varepsilon}) v w^{\varepsilon} \rangle_{H^{-1}(\Omega^{\varepsilon}), H^{1}_{0}(\Omega^{\varepsilon})} = \\ = \langle \mu^{\varepsilon}, S_{k,n}(\widetilde{u^{\varepsilon}}) v w^{\varepsilon} \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)} \to \\ \to \langle \mu, S_{k,n}(u^{0}) v \rangle_{H^{-1}(\Omega), H^{1}_{0}(\Omega)} = \int_{\Omega} S_{k,n}(u^{0}) v \, d\mu \quad \text{as } \varepsilon \to 0. \end{cases}$$

$$(8.25)$$

For what concerns the second and the third terms of the right-hand side of (8.23), we have, using Lebesgue's dominated convergence theorem as well as (6.5) and (8.24),

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$$-\int_{\Omega^{\varepsilon}} {}^{t}A(x)Dz^{\varepsilon}Dv \, S_{k,n}(u^{\varepsilon}) = -\int_{\Omega} {}^{t}A(x)Dz^{\varepsilon}Dv \, S_{k,n}(\widetilde{u^{\varepsilon}}) \to 0 \text{ as } \varepsilon \to 0, \quad (8.26)$$

$$\begin{cases} \int_{\Omega^{\varepsilon}} {}^{t}A(x)DvDS_{k,n}(u^{\varepsilon}) \, z^{\varepsilon} = \int_{\Omega} {}^{t}A(x)DvDS_{k,n}(\widetilde{u^{\varepsilon}}) \, z^{\varepsilon} \to \\ \to \int_{\Omega} {}^{t}A(x)DvDS_{k,n}(u^{0}) \quad \text{as } \varepsilon \to 0. \end{cases}$$

$$(8.27)$$

Collecting together (8.23), (8.25), (8.26) and (8.27), we have proved that the first term of the right-hand side of (8.22) satisfies, for n and k fixed with n > k > 0,

$$\begin{cases} \int_{\Omega^{\varepsilon}} {}^{t} A(x) D(z^{\varepsilon} v) DS_{k,n}(u^{\varepsilon}) \to \\ \to \int_{\Omega} {}^{t} A(x) Dv DS_{k,n}(u^{0}) + \int_{\Omega} S_{k,n}(u^{0}) v \, d\mu \quad \text{as } \varepsilon \to 0. \end{cases}$$

$$(8.28)$$

Let us now pass to the limit in the right-hand side of (8.28) as n tends to infinity.

Since $DS_{k,n}(u^0)=DG_k(u^0)\chi_{\{k\leq u^0\leq n\}},$ one has by Lebesgue's dominated convergence theorem

$$S_{k,n}(u^0) \to G_k(u^0)$$
 in $H_0^1(\Omega)$ strongly as $n \to +\infty$,

and therefore, still by Lebesgue's dominated convergence theorem,

$$S_{k,n}(u^0) \to G_k(u^0)$$
 in $L^1(\Omega; d\mu)$ strongly as $n \to +\infty$.

Therefore the right-hand side of (8.28) satisfies, since $v \in L^{\infty}(\Omega; d\mu)$ (see (2.26)),

$$\begin{cases} \int_{\Omega} {}^{t} A(x) Dv DS_{k,n}(u^{0}) + \int_{\Omega} S_{k,n}(u^{0}) v \, d\mu \to \\ \rightarrow \int_{\Omega} {}^{t} A(x) Dv DG_{k}(u^{0}) + \int_{\Omega} G_{k}(u^{0}) v \, d\mu \text{ as } n \to +\infty. \end{cases}$$

$$(8.29)$$

Passing to the limit in the right-hand side of (8.22), first as ε tends to zero for n fixed and then as n tends to infinity, and collecting together (8.28) and (8.29), will prove (8.17) whenever we will have proved that the second term of the right-hand side of (8.22) satisfies

$$\lim_{\varepsilon} \sup_{\Omega^{\varepsilon}} \left| \int_{\Omega^{\varepsilon}} {}^{t} A(x) D(z^{\varepsilon} v) DG_{n}(u^{\varepsilon}) \right| \to 0 \text{ as } n \to +\infty,$$
(8.30)

a result that we will prove in the third step just below.

Third step. In this third step we prove (8.30). As just said, this completes the proof of (8.17).

For that, we estimate the second term of the right-hand side of (8.22).

Since z^{ε} is bounded in $H^1(\Omega) \cap L^{\infty}(\Omega)$ (see (6.3) and (6.5)), one has

$$\left| \left| \int_{\Omega^{\varepsilon}} {}^{t} A(x) D(z^{\varepsilon}v) DG_{n}(u^{\varepsilon}) \right| \leq \\ \leq \|A\|_{(L^{\infty}(\Omega))^{N \times N}} \|D(z^{\varepsilon}v)\|_{(L^{2}(\Omega))^{N}} \|DG_{n}(u^{\varepsilon})\|_{(L^{2}(\Omega^{\varepsilon}))^{N}} \leq \\ \leq \|A\|_{(L^{\infty}(\Omega))^{N \times N}} \left(\|z^{\varepsilon}\|_{L^{\infty}(\Omega)} \|Dv\|_{(L^{2}(\Omega))^{N}} + \|v\|_{L^{\infty}(\Omega)} \|Dz^{\varepsilon}\|_{(L^{2}(\Omega))^{N}} \right) \\ \|DG_{n}(u^{\varepsilon})\|_{(L^{2}(\Omega^{\varepsilon}))^{N}} \quad \forall \varepsilon, \quad \forall n, \end{cases}$$

$$(8.31)$$

where C(v) is a constant which depends on v but neither on ε nor on n.

We now estimate $||DG_n(u^{\varepsilon})||_{(L^2(\Omega^{\varepsilon}))^N}$ in a way which is more precise than the a priori estimate (7.1). For that we use the (energy) equality (5.4) of [10], namely

$$\int_{\Omega^{\varepsilon}} A(x) DG_n(u^{\varepsilon}) DG_n(u^{\varepsilon}) = \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) G_n(u^{\varepsilon}),$$

which is formally obtained by using $G_n(u^{\varepsilon})$ as test function in (2.0^{ε}) . Using in this inequality the coercivity (2.2) of the matrix A and the growth condition (2.4) on the function F gives, since Γ is increasing and since $G_n(s) = 0$ for $s \leq n$,

$$\alpha \|DG_n(u^{\varepsilon})\|_{(L^2(\Omega^{\varepsilon}))^N}^2 \leq \int\limits_{\Omega^{\varepsilon}} \frac{h(x)}{\Gamma(u^{\varepsilon})} G_n(u^{\varepsilon}) \leq \int\limits_{\Omega^{\varepsilon}} \frac{h(x)}{\Gamma(n)} G_n(u^{\varepsilon}) = \int\limits_{\Omega} \frac{h(x)}{\Gamma(n)} G_n(\widetilde{u^{\varepsilon}}).$$

Passing to the limit in ε for *n* fixed thanks to (8.10) gives

$$\limsup_{\varepsilon} \|DG_n(u^{\varepsilon})\|^2_{(L^2(\Omega^{\varepsilon}))^N} \le \omega(n) \quad \forall n > 0,$$
(8.32)

where $\omega(n)$ is defined by

$$\omega^2(n) = \frac{1}{\alpha} \int_{\Omega} \frac{h(x)}{\Gamma(n)} G_n(u^0) \quad \forall n > 0.$$
(8.33)

Since Γ is increasing and since for $s \ge 0$ fixed $G_n(s)$ is nonincreasing in n, one has, for $n \ge n_0$,

$$\omega^{2}(n) = \frac{1}{\alpha} \int_{\Omega} \frac{h(x)}{\Gamma(n)} G_{n}(u^{0}) \chi_{\{u^{0} \ge n\}} \le \frac{1}{\alpha} \int_{\Omega} \frac{h(x)}{\Gamma(n_{0})} G_{n_{0}}(u^{0}) \chi_{\{u^{0} \ge n\}} \quad \forall n, n \ge n_{0}.$$

Since the measure of the set $\{x \in \Omega : u^0(x) \ge n\}$ tends to zero as n tends to infinity (recall that $u^0 \in L^2(\Omega)$ by (8.9)), and since $h(x)G_{n_0}(u^0) \in L^1(\Omega)$, one deduces, fixing n_0 , that

$$\omega^2(n) \to 0 \text{ as } n \to +\infty. \tag{8.34}$$

Collecting together (8.31), (8.32), (8.33) and (8.34) proves that the second term of the right-hand side of (8.22) satisfies

$$\limsup_{\varepsilon} \left| \int_{\Omega^{\varepsilon}} {}^{t} A(x) D(z^{\varepsilon} v) DG_{n}(u^{\varepsilon}) \right| \le C(v) \,\omega(n) \to 0, \text{ as } n \to +\infty.$$

i.e. (8.30), as announced in the beginning of this third step.

Fourth step. In this fourth step we pass to the limit, as ε tends to zero, in the second, third and fourth terms of the left-hand side of (8.16) and we prove that

$$\begin{cases} \sum_{i\in I} \int_{\Omega^{\varepsilon}} \hat{g}_i D(z^{\varepsilon} \hat{\varphi}_i T_k(u^{\varepsilon})) + \int_{\Omega^{\varepsilon}} G^{\varepsilon} D(w^{\varepsilon} v T_k(u^{\varepsilon})) + \\ + \int_{\Omega^{\varepsilon}} \left(z^{\varepsilon} \hat{f} - {}^t A(x) Dv Dz^{\varepsilon} - {}^t A(x) Dz^{\varepsilon} Dv \right) T_k(u^{\varepsilon}) \to \\ \to \sum_{i\in I} \int_{\Omega} \hat{g}_i D(\hat{\varphi}_i T_k(u^0)) + \int_{\Omega} \hat{f} T_k(u^0) + \int_{\Omega} T_k(u^0) v \, d\mu \quad \text{as } \varepsilon \to 0, \end{cases}$$

$$(8.35)$$

in which we observe that, in view of notation (3.3) of Definition 3.2, the two first terms of the right-hand side can also be written as

$$\sum_{i\in I} \int_{\Omega} \hat{g}_i D(\hat{\varphi}_i T_k(u^0)) + \int_{\Omega} \hat{f} T_k(u^0) = \langle \langle -div \, {}^t\!A(x) Dv, T_k(u^0) \rangle \rangle_{\Omega}.$$
(8.36)

For the second term of the left-hand side of (8.16), i.e. for the first term of the left-hand side of (8.35), we have, in view of (8.11) and since $\hat{\varphi}_i \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$,

$$\int_{\Omega^{\varepsilon}} \hat{g}_i D(z^{\varepsilon} \hat{\varphi}_i T_k(u^{\varepsilon})) = \int_{\Omega} \hat{g}_i D(z^{\varepsilon} \hat{\varphi}_i T_k(\widetilde{u^{\varepsilon}})) \to \int_{\Omega} \hat{g}_i D(\hat{\varphi}_i T_k(u^0)) \quad \text{as } \varepsilon \to 0.$$
(8.37)

Similarly, using the strong convergence of G^{ε} to G in $(L^2(\Omega))^N$ (see (6.10)), (8.13) and the fact that $v \in H^1_0(\Omega) \cap L^{\infty}(\Omega)$, we have, for the third term of the left-hand side of (8.16), i.e. for the second term of the left-hand side of (8.35),

$$\int_{\Omega^{\varepsilon}} G^{\varepsilon} D(w^{\varepsilon} v T_k(u^{\varepsilon})) = \int_{\Omega} G^{\varepsilon} D(w^{\varepsilon} v T_k(\widetilde{u^{\varepsilon}})) \to \int_{\Omega} G D(v T_k(u^0)) \quad \text{as } \varepsilon \to 0.$$
(8.38)

Moreover, in view of (6.10) and of (2.25) we have, since $vT_k(u^0) \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$ because of (8.11),

$$\begin{cases} \int_{\Omega} G D(vT_k(u^0)) = \langle -div \, G, vT_k(u^0) \rangle_{H^{-1}(\Omega), H^1_0(\Omega)} = \\ = \langle \mu, vT_k(u^0) \rangle_{H^{-1}(\Omega), H^1_0(\Omega)} = \int_{\Omega} T_k(u^0) \, v \, d\mu. \end{cases}$$
(8.39)

Finally, in view of (6.5) and (8.9), which implies that $T_k(\widetilde{u^{\varepsilon}})Dv$ converges to $T_k(u^0)Dv$ in $(L^2(\Omega))^N$ strongly by Lebesgue's dominated convergence theorem, we have, for the fourth term of the left-hand side of (8.16), i.e. for the third term of the left-hand side of (8.35),

$$\begin{cases} \int\limits_{\Omega^{\varepsilon}} \left(z^{\varepsilon} \hat{f} - {}^{t} A(x) Dv Dz^{\varepsilon} - {}^{t} A(x) Dz^{\varepsilon} Dv \right) T_{k}(u^{\varepsilon}) = \\ = \int\limits_{\Omega} \left(z^{\varepsilon} \hat{f} - {}^{t} A(x) Dv Dz^{\varepsilon} - {}^{t} A(x) Dz^{\varepsilon} Dv \right) T_{k}(\widetilde{u^{\varepsilon}}) \to \int\limits_{\Omega} \hat{f} T_{k}(u^{0}) \qquad (8.40) \\ \text{as } \varepsilon \to 0. \end{cases}$$

Collecting together (8.37), (8.38), (8.39) and (8.40) we have proved (8.35).

At this point, see (8.17), (8.18), (8.35) and (8.36), we passed to the limit in the left-hand side of (8.16). In the sixth, seventh and eighth steps below, we will pass to the limit in the right-hand side of (8.16).

Fifth step. Before of doing that, we prove in the present fifth step that

$$\int_{\Omega} F(x, u^0)v < +\infty \quad \forall v \in \mathcal{V}(\Omega), v \ge 0,$$
(8.41)

or in other terms that assertion (5.7 i) holds true.

Since the left-hand side of (8.16) converges as ε tends to zero, the right-hand side of (8.16) satisfies

$$\int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) z^{\varepsilon} v \le C(v) \quad \forall v \in \mathcal{V}(\Omega), v \ge 0, \ \forall \varepsilon,$$
(8.42)

where the constant $C(v) < +\infty$ does not depend on ε . Using the extension by zero defined in (2.12), (8.42) is equivalent to

$$\int_{\Omega} \widetilde{F(x, u^{\varepsilon})} z^{\varepsilon} v \le C(v) \ \forall \varepsilon.$$
(8.43)

We claim that

$$\widetilde{F(x, u^{\varepsilon})} \to F(x, u^0)$$
 a.e. $x \in \Omega$ as $\varepsilon \to 0$. (8.44)

Indeed, in view of Remark 6.5 and in particular of (6.27), we know that, for every subsequence ε' of ε and for almost every $x_0 \in \Omega$, there exists $\varepsilon_0(x_0)$ such that x_0 belongs to $\Omega^{\varepsilon'}$ for every $\varepsilon' < \varepsilon_0(x_0)$. This implies that

$$\widetilde{F(x_0, u^{\varepsilon'}(x_0))} = F(x_0, \widetilde{u^{\varepsilon'}}(x_0)) \quad \forall \varepsilon', \, \varepsilon' < \varepsilon_0(x_0).$$

Since

$$F(x_0, \widetilde{u^{\varepsilon'}}(x_0)) \to F(x_0, u^0(x_0))$$
 a.e. $x_0 \in \Omega$

in view of the convergence (8.9) and of the Carathéodory hypothesis (2.3), this implies (8.44).

Results (8.43) and (8.44) combined with (6.5), the fact that $F(x, \widetilde{u^{\varepsilon}})z^{\varepsilon}v \geq 0$, and finally Fatou's lemma immediately imply (8.41).

Sixth step. As already said at the end of the fourth step, in the present sixth step and in the seventh and eighth steps, we pass to the limit in the right-hand side of (8.16).

For doing that, we introduce a new parameter $\delta > 0$ and we write the right-hand side of (8.16) as

$$\int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) z^{\varepsilon} v = \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) Z_{\delta}(u^{\varepsilon}) z^{\varepsilon} v + \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) \left(1 - Z_{\delta}(u^{\varepsilon})\right) z^{\varepsilon} v, \quad (8.45)$$

where Z_{δ} is the function defined by (7.6).

In the present sixth step we prove that the first term of the right-hand side of (8.45) satisfies

$$\limsup_{\varepsilon} \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) Z_{\delta}(u^{\varepsilon}) z^{\varepsilon} v \to 0 \quad \text{as } \delta \to 0.$$
(8.46)

For that we use estimate (7.8) of Proposition 7.5 above with $v^{\varepsilon} = z^{\varepsilon}v$ for any $v \in \mathcal{V}(\Omega)$, $v \geq 0$; this choice is licit in view of (6.4). In view of (6.11), the estimate (7.8) reads as

$$\begin{cases} \forall \delta > 0, \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) Z_{\delta}(u^{\varepsilon}) z^{\varepsilon} v \leq I_{\delta}^{\varepsilon} + II_{\delta}^{\varepsilon}, \\ \text{where} \\ I_{\delta}^{\varepsilon} = \frac{3}{2} \left(\int_{\Omega^{\varepsilon}} \left| \sum_{i \in I} \hat{g}_{i} D(z^{\varepsilon} \hat{\varphi}_{i}) + G^{\varepsilon} D(w^{\varepsilon} v) + z^{\varepsilon} \hat{f} - {}^{t} A(x) Dv Dz^{\varepsilon} - {}^{t} A(x) Dz^{\varepsilon} Dv \right| \right) \delta, \\ II_{\delta}^{\varepsilon} = \sum_{i \in I} \int_{\Omega^{\varepsilon}} Z_{\delta}(u^{\varepsilon}) \hat{g}_{i} Du^{\varepsilon} z^{\varepsilon} \hat{\varphi}_{i} + \int_{\Omega^{\varepsilon}} Z_{\delta}(u^{\varepsilon}) G^{\varepsilon} Du^{\varepsilon} w^{\varepsilon} v. \end{cases}$$

$$(8.47)$$

Since z^{ε} and w^{ε} are bounded in $H^1(\Omega) \cap L^{\infty}(\Omega)$ (see (6.5) and (2.17)), since $\hat{\varphi}_i$ and v belong to $H^1_0(\Omega) \cap L^{\infty}(\Omega)$, since \hat{g}_i belongs to $(L^2(\Omega))^N$, since G^{ε} is bounded in $(L^2(\Omega))^N$ (see (6.10)), and since \hat{f} belongs to $L^1(\Omega)$, we have, as far as I^{ε}_{δ} is concerned,

$$\limsup_{\varepsilon} I_{\delta}^{\varepsilon} \le \frac{3}{2} C\delta, \tag{8.48}$$

where the constant C does not depend neither on ε nor on δ , and therefore we have

$$\limsup_{\varepsilon} I^{\varepsilon}_{\delta} \to 0 \quad \text{as } \delta \to 0.$$
(8.49)

For what concerns the first term of $II_{\delta}^{\varepsilon}$, in view of (6.8), (3.9), (3.7^{ε} i), (3.7^{ε} iii) and (3.7^{ε} iv), we have

$$\begin{cases} Du^{\varepsilon} z^{\varepsilon} \hat{\varphi}_{i} = (DT_{k}(u^{\varepsilon}) + DG_{k}(u^{\varepsilon})) z^{\varepsilon} \hat{\varphi}_{i} = \\ = D(z^{\varepsilon} \hat{\varphi}_{i} T_{k}(u^{\varepsilon})) - T_{k}(u^{\varepsilon}) D(z^{\varepsilon} \hat{\varphi}_{i}) + DG_{k}(u^{\varepsilon}) z^{\varepsilon} \hat{\varphi}_{i} & \text{in } \mathcal{D}'(\Omega^{\varepsilon}). \end{cases}$$

Since when $a^{\varepsilon} \in L^2(\Omega^{\varepsilon})$ and $b^{\varepsilon} \in L^2(\Omega)$, one has $\widetilde{a^{\varepsilon}b^{\varepsilon}} = \widetilde{a^{\varepsilon}b^{\varepsilon}}$, and since $z^{\varepsilon}\hat{\varphi}_i T_k(u^{\varepsilon})$ and $G_k(u^{\varepsilon})$ belong to $H_0^1(\Omega^{\varepsilon})$, we have in view of (2.13)

$$\begin{cases} \widetilde{Du^{\varepsilon} z^{\varepsilon} \hat{\varphi}_{i}} = \widetilde{D(z^{\varepsilon} \hat{\varphi}_{i} T_{k}(u^{\varepsilon}))} - \widetilde{T_{k}(u^{\varepsilon})} D(z^{\varepsilon} \varphi_{i}) + \widetilde{DG_{k}(u^{\varepsilon})} z^{\varepsilon} \hat{\varphi}_{i} = \\ = D(\widetilde{z^{\varepsilon} \hat{\varphi}_{i} T_{k}(u^{\varepsilon})}) - \widetilde{T_{k}(u^{\varepsilon})} D(z^{\varepsilon} \hat{\varphi}_{i}) + D\widetilde{G_{k}(u^{\varepsilon})} z^{\varepsilon} \hat{\varphi}_{i} \quad \text{in } \mathcal{D}'(\Omega). \end{cases}$$

But since $T_k(0) = G_k(0) = 0$ implies that $\widetilde{T_k(u^{\varepsilon})} = T_k(\widetilde{u^{\varepsilon}})$ and $\widetilde{G_k(u^{\varepsilon})} = G_k(\widetilde{u^{\varepsilon}})$, we finally have

$$\widetilde{Du^{\varepsilon}z^{\varepsilon}\hat{\varphi_{i}}} = D(z^{\varepsilon}\hat{\varphi_{i}}T_{k}(\widetilde{u^{\varepsilon}})) - T_{k}(\widetilde{u^{\varepsilon}})D(z^{\varepsilon}\hat{\varphi_{i}}) + DG_{k}(\widetilde{u^{\varepsilon}})z^{\varepsilon}\hat{\varphi_{i}} \text{ in } \mathcal{D}'(\Omega).$$

Since $\widetilde{Du^{\varepsilon}z^{\varepsilon}\hat{\varphi}_i}$ is zero on $\Omega \setminus \Omega^{\varepsilon}$, and even if

$$\widetilde{Z_{\delta}(u^{\varepsilon})} = 0$$
 in $\Omega \setminus \Omega^{\varepsilon}$ while $Z_{\delta}(\widetilde{u^{\varepsilon}}) = 1$ in $\Omega \setminus \Omega^{\varepsilon}$,

we have

$$\begin{cases} \widetilde{Z_{\delta}(u^{\varepsilon})\hat{g}_{i}Du^{\varepsilon}z^{\varepsilon}\varphi_{i}} = Z_{\delta}(\widetilde{u^{\varepsilon}})\hat{g}_{i}\widetilde{Du^{\varepsilon}z^{\varepsilon}\varphi_{i}} = \\ = Z_{\delta}(\widetilde{u^{\varepsilon}})\hat{g}_{i}\left(D(z^{\varepsilon}\hat{\varphi}_{i}T_{k}(\widetilde{u^{\varepsilon}})) - T_{k}(\widetilde{u^{\varepsilon}})D(z^{\varepsilon}\hat{\varphi}_{i}) + DG_{k}(\widetilde{u^{\varepsilon}})z^{\varepsilon}\hat{\varphi}_{i}\right) \text{ in } \mathcal{D}'(\Omega). \end{cases}$$

$$(8.50)$$

Therefore, in view of (8.9), (8.10), (8.11) and (6.5), we have

$$\begin{cases} \forall \delta > 0 \text{ fixed,} \\ \int Z_{\delta}(u^{\varepsilon}) \, \hat{g}_{i} \, Du^{\varepsilon} \, z^{\varepsilon} \hat{\varphi}_{i} = \\ = \int Z_{\delta}(\widetilde{u^{\varepsilon}}) \, \hat{g}_{i} \left(D(z^{\varepsilon} \hat{\varphi}_{i} T_{k}(\widetilde{u^{\varepsilon}})) - T_{k}(\widetilde{u^{\varepsilon}}) D(z^{\varepsilon} \hat{\varphi}_{i}) + DG_{k}(\widetilde{u^{\varepsilon}}) z^{\varepsilon} \hat{\varphi}_{i} \right) \rightarrow \\ \to \int Z_{\delta}(u^{0}) \, \hat{g}_{i} \left(D(\hat{\varphi}_{i} T_{k}(u^{0})) - T_{k}(u^{0}) D\hat{\varphi}_{i} + DG(u^{0}) \hat{\varphi}_{i} \right) = \\ = \int Z_{\delta}(u^{0}) \, \hat{g}_{i} \, Du^{0} \, \hat{\varphi}_{i} \quad \text{as } \varepsilon \rightarrow 0. \end{cases}$$

$$(8.51)$$

A proof which is very similar to the proof of (8.51) implies that for the second term of $II_{\delta}^{\varepsilon}$ we have

$$\begin{cases} \forall \delta > 0 \text{ fixed,} \\ \int_{\Omega^{\varepsilon}} Z_{\delta}(u^{\varepsilon}) G^{\varepsilon} D u^{\varepsilon} w^{\varepsilon} v \to \int_{\Omega} Z_{\delta}(u^{0}) G^{0} D u^{0} v \quad \text{as } \varepsilon \to 0. \end{cases}$$

$$(8.52)$$

Let us now pass to the limit in the right-hand side of (8.51) as δ tends to zero. Since

$$Z_{\delta}(s) \to \chi_{\{s=0\}}(s), \quad \forall s \ge 0,$$

and since $u^0 \in H^1_{\text{loc}}(\Omega)$ implies that

$$Du^0 = 0$$
 a.e. in $\{x \in \Omega : u^0(x) = 0\}$,

we have

$$\int_{\Omega} Z_{\delta}(u^0) \,\hat{g}_i \, Du^0 \hat{\varphi}_i \to \int_{\Omega} \chi_{\{u^0=0\}} \hat{g}_i \, Du^0 \,\hat{\varphi}_i = 0 \quad \text{as } \delta \to 0.$$
(8.53)

The same proof implies that for the right-hand side of (8.52) we have

$$\int_{\Omega} Z_{\delta}(u^0) G^0 D u^0 v \to 0 \quad \text{as } \delta \to 0.$$
(8.54)

Collecting together the definition (8.47) of $II_{\delta}^{\varepsilon}$ and the results obtained in (8.51), (8.53), (8.52) and (8.54) proves that

$$\lim_{\varepsilon} II^{\varepsilon}_{\delta} \to 0 \quad \text{as } \delta \to 0.$$
(8.55)

Finally, collecting together (8.47), (8.49) and (8.55) proves (8.46), as announced at the beginning of this sixth step.

Seventh step. In this seventh step we prove that

$$\int_{\Omega \cap \{u^0 = 0\}} F(x, u^0) v = 0.$$
(8.56)

By (2.12), we have $\widetilde{F(x,u^{\varepsilon})} = 0$ in $\Omega \setminus \Omega^{\varepsilon}$, and even if $Z_{\delta}(\widetilde{u^{\varepsilon}}) = 1$ in $\Omega \setminus \Omega^{\varepsilon}$, we have

$$\widetilde{F(x,u^{\varepsilon})Z_{\delta}(u^{\varepsilon})} = \widetilde{F(x,u^{\varepsilon})}Z_{\delta}(\widetilde{u^{\varepsilon}}) \text{ in } \Omega, \qquad (8.57)$$

from which we deduce that

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$$\begin{cases} \forall \delta > 0, \int\limits_{\Omega \cap \{u^0 = 0\}} \widetilde{F(x, u^{\varepsilon})} Z_{\delta}(\widetilde{u^{\varepsilon}}) z^{\varepsilon} v = \int\limits_{\Omega \cap \{u^0 = 0\}} \widetilde{F(x, u^{\varepsilon})} Z_{\delta}(u^{\varepsilon}) z^{\varepsilon} v = \\ = \int\limits_{\Omega^{\varepsilon} \cap \{u^0 = 0\}} F(x, u^{\varepsilon}) Z_{\delta}(u^{\varepsilon}) z^{\varepsilon} v \leq \int\limits_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) Z_{\delta}(u^{\varepsilon}) z^{\varepsilon} v. \end{cases}$$
(8.58)

On the other hand, (8.44), (8.9) and (6.5) imply that

$$\widetilde{F(x,u^{\varepsilon})}Z_{\delta}(\widetilde{u^{\varepsilon}})z^{\varepsilon}v \to F(x,u^{0})Z_{\delta}(u^{0})v \text{ a.e. in } \Omega \text{ as } \varepsilon \to 0.$$
(8.59)

Using Fatou's lemma in the left-hand side of (8.58) and observing that

$$Z_{\delta}(u^0) = 1$$
 in $\{x \in \Omega : u^0(x) = 0\} \ \forall \delta > 0,$

the almost everywhere convergence (8.59) and the inequality (8.58) imply that

$$\begin{cases} \forall \delta > 0, \int_{\Omega \cap \{u^0 = 0\}} F(x, u^0) v \leq \liminf_{\varepsilon} \int_{\Omega \cap \{u^0 = 0\}} \widetilde{F(x, u^{\varepsilon})} Z_{\delta}(\widetilde{u^{\varepsilon}}) z^{\varepsilon} v \leq \\ \leq \limsup_{\varepsilon} \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) Z_{\delta}(u^{\varepsilon}) z^{\varepsilon} v, \end{cases}$$

$$(8.60)$$

which letting δ tend to zero and using (8.46) implies (8.56).

Eighth step. In this eighth step we prove that the second term of the right-hand side of (8.45) satisfies

$$\lim_{\varepsilon} \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon}) (1 - Z_{\delta}(u^{\varepsilon})) \, z^{\varepsilon} v \to \int_{\Omega} F(x, u^{0}) \, v \text{ as } \delta \to 0.$$
(8.61)

Indeed, similarly to the results obtained in the seventh step, we have

$$\widetilde{F(x,u^{\varepsilon})(1-Z_{\delta}(u^{\varepsilon}))} = \widetilde{F(x,u^{\varepsilon})}(1-Z_{\delta}(\widetilde{u^{\varepsilon}})) \text{ in } \Omega,$$
(8.62)

as well as

$$\widetilde{F(x,u^{\varepsilon})}(1-Z_{\delta}(\widetilde{u^{\varepsilon}}))z^{\varepsilon}v \to F(x,u^{0})(1-Z_{\delta}(u^{0}))v \text{ a.e. in }\Omega.$$
(8.63)

On the other hand, we have

$$1 - Z_{\delta}(\widetilde{u^{\varepsilon}}) = 0 \text{ a.e. in } \{ x \in \Omega : \widetilde{u^{\varepsilon}}(x) \le \delta \},\$$

while the fact that the set $\{x \in \Omega : \widetilde{u^{\varepsilon}}(x) > \delta\}$ is a subset of Ω^{ε} (since $\widetilde{u^{\varepsilon}} = 0$ in $\Omega \setminus \Omega^{\varepsilon}$) and the conditions (2.4 *iii*) and (2.4 *iii*) on the functions F(x, s) and $\Gamma(s)$ imply that

$$0 \leq \widetilde{F(x,u^{\varepsilon})} = F(x,u^{\varepsilon}) \leq \frac{h(x)}{\Gamma(u^{\varepsilon})} = \frac{h(x)}{\Gamma(\widetilde{u^{\varepsilon}})} \leq \frac{h(x)}{\Gamma(\delta)} \text{ a.e. in } \{x \in \Omega : \widetilde{u^{\varepsilon}}(x) > \delta\}.$$

Together with $0 \le Z_{\delta}(s) \le 1$ and (6.3), the two latest assertions imply that

$$0 \le \widetilde{F(x, u^{\varepsilon})} (1 - Z_{\delta}(\widetilde{u^{\varepsilon}})) z^{\varepsilon} v \le \frac{h(x)}{\Gamma(\delta)} v \text{ a.e. in } \Omega,$$
(8.64)

where $h(x)v \in L^1(\Omega)$ in view of condition (2.4 *i*).

From (8.62), (8.63), (8.64) and Lebesgue's dominated convergence theorem we deduce that

$$\begin{cases} \forall \delta > 0, \ \int_{\Omega^{\varepsilon}} F(x, u^{\varepsilon})(1 - Z_{\delta}(u^{\varepsilon})) \ z^{\varepsilon}v = \int_{\Omega} \widetilde{F(x, u^{\varepsilon})} \left(1 - Z_{\delta}(\widetilde{u^{\varepsilon}})\right) z^{\varepsilon}v \rightarrow \\ \rightarrow \int_{\Omega} F(x, u^{0}) \left(1 - Z_{\delta}(u^{0})\right) v \text{ as } \varepsilon \rightarrow 0. \end{cases}$$

$$(8.65)$$

Since

$$Z_{\delta}(u^0) \to \chi_{\{u^0=0\}}$$
 a.e. in Ω as $\delta \to 0$,

applying again Lebesgue's dominated convergence theorem thanks to (8.41) and then using (8.56) implies that

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$$\int_{\Omega} F(x, u^0)(1 - Z_{\delta}(u^0)) v \to \int_{\Omega} F(x, u^0)(1 - \chi_{\{u^0 = 0\}}) v = \int_{\Omega} F(x, u^0) v \text{ as } \delta \to 0.$$

This proves (8.61).

Ninth (and last) step. Collecting together (8.16), (8.17), (8.18), (8.35), (8.36), (8.45), (8.46) and (8.61), we have proved that u^0 satisfies (5.7 *ii*).

Recall that we have also proved in the fifth step above (see (8.41)) that u^0 satisfies (5.7 *i*).

Finally recall that we have proved in the first step above that the subsequence that we have extracted satisfies convergences (5.3), (5.4) and (5.5), and that its limit u^0 satisfies (5.6).

The proof of Theorem 5.1 is complete. \Box

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